

Group-invariant solutions of semilinear Schrödinger equations in multi-dimensions

Stephen C. Anco and Wei Feng

Citation: [Journal of Mathematical Physics](#) **54**, 121504 (2013); doi: 10.1063/1.4830316

View online: <http://dx.doi.org/10.1063/1.4830316>

View Table of Contents: <http://scitation.aip.org/content/aip/journal/jmp/54/12?ver=pdfcov>

Published by the [AIP Publishing](#)

Articles you may be interested in

[Variable coefficient nonlinear Schrödinger equations with four-dimensional symmetry groups and analysis of their solutions](#)

J. Math. Phys. **52**, 093702 (2011); 10.1063/1.3634005

[q -conformal invariant equations and q -plane wave solutions](#)

J. Math. Phys. **45**, 3788 (2004); 10.1063/1.1790049

[Group theory approach to the Dirac equation with a Coulomb plus scalar potential in D+1 dimensions](#)

J. Math. Phys. **44**, 4467 (2003); 10.1063/1.1604185

[The Ermanno–Bernoulli constants and representations of the complete symmetry group of the Kepler problem](#)

J. Math. Phys. **44**, 4090 (2003); 10.1063/1.1576903

[Lie groups and differential equations: symmetries, conservation laws, and exact solutions of mathematical models in physics](#)

Phys. Part. Nucl. **28**, 241 (1997); 10.1134/1.953039



Group-invariant solutions of semilinear Schrödinger equations in multi-dimensions

Stephen C. Anco^{1,a)} and Wei Feng^{1,2,b)}

¹*Department of Mathematics, Brock University, St. Catharines, Ontario L2S3A1, Canada*

²*Department of Mathematics, Zhejiang University of Technology, Hangzhou 310014, China*

(Received 21 January 2013; accepted 31 October 2013; published online 19 December 2013)

Symmetry group methods are applied to obtain all explicit group-invariant radial solutions to a class of semilinear Schrödinger equations in dimensions $n \neq 1$. Both focusing and defocusing cases of a power nonlinearity are considered, including the special case of the pseudo-conformal power $p = 4/n$ relevant for critical dynamics. The methods involve, first, reduction of the Schrödinger equations to group-invariant semilinear complex 2nd order ordinary differential equations (ODEs) with respect to an optimal set of one-dimensional point symmetry groups, and second, use of inherited symmetries, hidden symmetries, and conditional symmetries to solve each ODE by quadratures. Through Noether's theorem, all conservation laws arising from these point symmetry groups are listed. Some group-invariant solutions are found to exist for values of n other than just positive integers, and in such cases an alternative two-dimensional form of the Schrödinger equations involving an extra modulation term with a parameter $m = 2 - n \neq 0$ is discussed. © 2013 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4830316>]

I. INTRODUCTION

Semilinear Schrödinger equations with power nonlinearities for $u(t, x)$, $x \in \mathbb{R}^n$,

$$iu_t = \Delta u + k|u|^p u, \quad p \neq 0, \quad k = \text{const} \neq 0, \quad (1.1)$$

provide models of many interesting physical phenomena,¹ such as propagation of laser beams in nonlinear media, slow oscillations of plasma waves, motion of water waves at the free surface of an ideal fluid, dynamics of imperfect Bose condensates, and continuous-limits for mesoscopic molecular structure. In multi-dimensions $n > 1$, solutions exhibit very rich types of behaviour,^{1,2} particularly radial similarity solutions and more general group-invariant radial solutions which are important for investigating collapse or blow-up behaviour, dispersive behaviour, critical dynamics, and asymptotic attractors, as well as for testing numerical solution methods.

To date in the literature, only a few explicit n -dimensional radial solutions $u(t, |x|)$ are apparently known (e.g., see Ref. 3). Most of the work on explicit solutions to Schrödinger equations with power nonlinearities (1.1) has concentrated on systematically applying symmetry methods^{4,5} to classify all possible types of group-invariant solutions,^{6–9} including radial and cylindrical as well as other less geometric types, only for the pseudo-conformal power $p = 4/3$ in dimension $n = 3$ in addition to the lowest even powers $p = 2, 4$ in dimensions $n = 2, 3$. These are the cases of greatest relevance for physical applications.

The present paper, in contrast, will be devoted to deriving all group-invariant radial solutions $u(t, |x|)$ of the Schrödinger equation (1.1) in all dimensions $n \neq 1$ and for all powers $p \neq 0$, including the case of the general pseudo-conformal power $p = 4/n$ relevant for understanding critical dynamics in arbitrary dimensions. Both the focusing case $k > 0$ and the defocusing case $k < 0$ will be considered.

a) sanco@brocku.ca

b) fengwei9999@gmail.com

Radial solutions $u(t, |x|)$ of the Schrödinger equation (1.1) satisfy the corresponding radial equation

$$iu_t = u_{rr} + (n-1)r^{-1}u_r + k|u|^p u, \quad n \neq 1 \quad (1.2)$$

with $r = |x|$. Equation (1.2) describes a general model for the slow modulation of n -dimensional radial waves in weakly nonlinear, dispersive, isotropic media. In particular, the amplitude for harmonic waves $\exp(i(\kappa r - \omega t))$ with a dispersion relation $\omega = \omega(\kappa)$ in such a medium is given by¹ $u(\epsilon t, \epsilon(r - \omega_\kappa t))$ in terms of a small parameter $|\epsilon| \ll 1$.

Group-invariant solutions $u(t, r)$ of Eq. (1.2) arise from reductions by one-dimensional groups of point transformations on r, t, u, \bar{u} that leave the equation invariant. Each such symmetry reduction yields a complex 2nd order semilinear ordinary differential equation (ODE)

$$U'' = f(\xi, U, \bar{U}, U', \bar{U}') \quad (1.3)$$

formulated in terms of the invariants

$$\xi = \Xi(t, r, u, \bar{u}), \quad U = \Upsilon(t, r, u, \bar{u}), \quad \bar{U} = \bar{\Upsilon}(t, r, u, \bar{u}) \quad (1.4)$$

determined by a given symmetry transformation group, provided that this system (1.4) can be inverted (at least implicitly) to obtain both of the dependent variables u, \bar{u} , and one of the independent variables r, t in terms of ξ, U, \bar{U} , and the other independent variable (which will be the case⁴ whenever the orbits of the transformation group acting on the variables (r, t, u, \bar{u}) are one-dimensional and have a projectable regular action on (r, t)). Then each solution of the ODE (1.3) for $U(\xi)$ will yield a group-invariant solution of the radial Schrödinger equation (1.2) for $u(t, r)$.

If two groups of point symmetry transformations are related by conjugation with respect to some point transformation in the full symmetry group of the radial Schrödinger equation (1.2), then the action of this point transformation on solutions $u(t, r)$ will map the group-invariant solutions determined by the two symmetry groups into each other. Consequently, for the purpose of finding all group-invariant radial solutions, it is sufficient to work with a maximal set of one-dimensional point symmetry groups that are conjugacy inequivalent. For each such group, once the resulting group-invariant solutions have been found, the full symmetry group of the radial Schrödinger equation can be applied on these solutions to obtain the group-invariant solutions determined by all other one-dimensional point symmetry groups in the same conjugacy class.

Interestingly, group-invariant solutions $u(t, r)$ exist for values of n other than just positive integers. In such cases Eq. (1.2) can be interpreted alternatively as modeling the slow modulation of 2-dimensional radial waves in a planar, weakly nonlinear, dispersive, isotropic medium with a point-source disturbance at the origin. Specifically, this equation can be written in the equivalent form

$$iu_t = u_{rr} + (1-m)r^{-1}u_r + k|u|^p u \quad (1.5)$$

with a parameter $m = 2 - n$ which is applicable for any value of n . The term $mr^{-1}u_r$ has a natural interpretation through the net modulation defined by the 2-dimensional integral

$$\mathcal{C}(t) = \int_0^\infty u \, r \, dr \quad (1.6)$$

on the radial domain. If u is regular at the origin (i.e., $\lim_{r \rightarrow 0} u_r = 0$), then this integral formally satisfies

$$\frac{d\mathcal{C}}{dt} = -(S + m \lim_{r \rightarrow 0} iu), \quad (1.7)$$

where

$$S(t) = ik \int_0^\infty u|u|^p \, r \, dr \quad (1.8)$$

measures the net amount of modulation produced by the nonlinear term in Eq. (1.5). Note that \mathcal{C} is conserved, $\frac{d\mathcal{C}}{dt} = 0$, whenever both the modulation term and the nonlinear term are absent,

$m = k = 0$ (describing a linear, non-dispersive, planar medium). Thus, when $m \neq 0$, the modulation term $mr^{-1}u_r$ models the effect of a point-source disturbance at the origin $r = 0$, which alters the slow modulation of harmonic waves in a planar, weakly nonlinear, dispersive medium.

For a given point symmetry reduction of the radial Schrödinger equation (1.2), or the equivalent 2-dimensional equation (1.5), the task of solving the complex 2nd order semilinear ODE (1.3) for $U(\xi)$ is typically very non-trivial. One systematic approach is the method of reduction of order,^{4,5} which relies essentially on the algebraic structure of the group of point symmetries admitted by this ODE and on whether the ODE admits a Lagrangian structure or not, depending on the values of the parameters p and n (or $m = 2 - n$). Reduction of order is simplest to carry out using an equivalent polar system of 2nd order real ODEs for the amplitude and phase of the variable

$$U = A \exp(i\Phi). \quad (1.9)$$

First, Eq. (1.2) together with its complex-conjugate equation are the respective Euler–Lagrange equations $\delta L/\delta \bar{u} = 0$ and $\delta L/\delta u = 0$ of the radial Schrödinger Lagrangian

$$L = (\frac{1}{2}i\bar{u}u_t - \frac{1}{2}i\bar{u}_t u + \bar{u}_r u_r - \frac{2}{p+2}k\bar{u}^{1+p/2}u^{1+p/2})r^{n-1}. \quad (1.10)$$

As a consequence, if this Lagrangian is invariant under the given point symmetry transformations used for the reduction, then the polar system for $A(\xi)$ and $\Phi(\xi)$ will inherit a Lagrangian obtained by reduction of the radial Euler–Lagrange structure. Second, the phase rotation transformations

$$u \rightarrow e^{i\phi}u, \quad \bar{u} \rightarrow e^{-i\phi}\bar{u} \quad (1.11)$$

comprise a U(1) group of point symmetries of Eq. (1.2). This group is readily shown to commute with the full group of point symmetries admitted by Eq. (1.2). Therefore, if the given point symmetry reduction starts from any other one-dimensional group of point symmetries, then the polar system for $A(\xi)$ and $\Phi(\xi)$ will inherit a U(1) group of point symmetry transformations on ξ , A , Φ obtained by reduction of the phase rotation transformations on u , \bar{u} . Moreover, this system will also inherit all point symmetries that belong to the normalizer subgroup of the given point symmetry used for the reduction. An important remark is that the full group of point symmetries admitted by the system can possibly contain “hidden” point symmetries in addition to those point symmetries inherited via reduction.

In the case when a given point symmetry reduction of the radial Schrödinger equation (1.2) inherits a Lagrangian structure, the polar system for $A(\xi)$ and $\Phi(\xi)$ can be solved explicitly (up to quadratures) if its Lagrangian is invariant under a two-dimensional group of point symmetry transformations. In the alternative case when no Lagrangian structure exists, a four-dimensional point symmetry group with a solvable Lie algebra structure is needed to obtain the explicit solutions (up to quadratures) for $A(\xi)$ and $\Phi(\xi)$. In all cases, each one-dimensional group of point symmetry transformations admitted by the polar system for $A(\xi)$ and $\Phi(\xi)$, or equivalently by the complex semilinear ODE for $U(\xi)$, can be used to find a single solution that is invariant under the admitted symmetry transformations on ξ , A , Φ .

A different reduction method is applicable in the Lagrangian cases. If a given polar system for $A(\xi)$ and $\Phi(\xi)$ inherits a U(1)-invariant Lagrangian obtained by reduction of both the radial Schrödinger Lagrangian (1.10) and the U(1) group of phase rotation symmetries (1.11), then the 2nd order ODE for $\Phi(\xi)$ can be explicitly integrated so that the system reduces to a 2nd order ODE for $A(\xi)$ alone, with $\Phi(\xi)$ given by an integral in terms of $A(\xi)$ containing an arbitrary constant. This reduced ODE for $A(\xi)$ has two useful features. First, all nonlinear terms in the reduced ODE vanish if the nonlinearity power is $p = -1$ and the arbitrary constant in integral for $\Phi(\xi)$ is zero. In this case, the resulting linear ODE for $A(\xi)$ can be solved explicitly in terms of special functions. Second, a Lagrangian structure can be derived for the reduced ODE starting from the U(1)-invariant Lagrangian of the original polar system. In the nonlinear case this ODE can be solved (up to quadratures) if its Lagrangian is invariant under a one-dimensional group of point symmetry transformations on ξ , A . In all cases the reduced ODE also can be solved (up to quadratures) if it admits a two-dimensional group of non-variational point symmetry transformations. More generally, any one-dimensional group of point symmetries admitted by the reduced ODE can be used to find a single invariant solution for $A(\xi)$. Each solution found for $A(\xi)$ determines a solution for $\Phi(\xi)$ and hence yields a solution to the

polar system. These solutions will turn out to differ from those solutions obtained via the previous reduction method if the group of point symmetries used to obtain $A(\xi)$ is a “hidden” group which is not inherited under reduction from any point symmetries admitted by the polar system for $A(\xi)$ and $\Phi(\xi)$. In particular, a “hidden” symmetry group can arise from point transformations on ξ , A that leave invariant the integral for $\Phi(\xi)$.

The rest of this paper is organized as follows.

Section II contains some preliminaries on symmetries and Noether’s theorem. We first state the full point symmetry structure for both the n -dimensional Schrödinger equation (1.1) and the radial Schrödinger equation (1.2). In addition, we summarize the conservation laws that arise from the Lagrangian structure of these two equations via Noether’s theorem. Next for the radial Schrödinger equation (1.2) we present a maximal set of one-dimensional point symmetry groups that are conjugacy inequivalent. Modulo phase rotations, the symmetry groups in this set consist of time translations, scalings, and a particular combination of inversions (pseudo-conformal transformations) and time translations. These three symmetry groups will be referred to as the optimal subgroups for symmetry considerations. We then write down the complex 2nd order semilinear ODEs (1.3) given by reduction under each point symmetry subgroup in the optimal set, and we summarize the Lagrangian structure admitted by each of the ODEs, depending on p and n .

Section III explains the reduction of order method in detail for solving complex 2nd order semilinear ODEs (1.3) by use of point symmetries. In particular, we show how to streamline the standard reduction steps in an efficient way by combining the Lagrangian and non-Lagrangian cases through the use of canonical coordinates determined by any admitted one-dimensional point symmetry group. We also explain details of the alternative reduction method in the $U(1)$ -invariant Lagrangian cases, which we carry out by again using canonical coordinates to streamline the steps.

Sections IV–VI apply these reduction methods to derive solutions $U(\xi)$ for each ODE (1.3) arising from the three optimal subgroups of point symmetries for the radial Schrödinger equation (1.2). In particular, we are able to obtain explicit solutions in terms of elementary functions.

The main results of the paper are presented in Sec. VII. We first list all of the group-invariant radial solutions $u(t, r)$ determined by applying the full group of point symmetries of the radial Schrödinger equation (1.2) to each of the group-invariant solutions $U(\xi)$ derived from the three optimal subgroups of point symmetries. We next discuss a few analytical features of these solutions $u(t, r)$, including cases that involve non-integer values of n .

Finally, Sec. VIII has some concluding remarks as well as comments on group-invariant radial solutions pertaining to blow-up behaviour of $u(t, r)$ for powers $p \geq 4/n$.

II. SYMMETRIES AND CONSERVATION LAWS

For the Schrödinger equation (1.1) in \mathbb{R}^n , a *point symmetry* is a one-dimensional Lie group of transformations acting on the variables (t, x, u, \bar{u}) such that the prolongation $\text{pr}X$ of its infinitesimal generator

$$X = \tau(t, x, u, \bar{u})\partial/\partial t + \zeta(t, x, u, \bar{u}) \cdot \partial/\partial x + \eta(t, x, u, \bar{u})\partial/\partial u + \bar{\eta}(t, x, u, \bar{u})\partial/\partial \bar{u} \quad (2.1)$$

satisfies $\text{pr}X(iu_t - \Delta u - k|u|^p u) = 0$ for all formal solutions $u(t, x)$ of Eq. (1.1). This is the condition for Eq. (1.1) to be infinitesimally invariant under the transformation group generated by X . Each such generator acting on solutions $u(t, x)$ has an equivalent *characteristic form*

$$\hat{X} = P(t, x, u, \bar{u}, u_t, \bar{u}_t, \nabla u, \nabla \bar{u})\partial/\partial u + \bar{P}(t, x, u, \bar{u}, u_t, \bar{u}_t, \nabla u, \nabla \bar{u})\partial/\partial \bar{u} \quad (2.2)$$

with

$$P = \eta - \tau u_t - \zeta \cdot \nabla u \quad (2.3)$$

satisfying

$$iD_t P - D_x \cdot D_x P - k(1 + p/2)\bar{u}^{p/2}u^{p/2}P - k(p/2)\bar{u}^{p/2-1}u^{1+p/2}\bar{P} = 0 \quad (2.4)$$

as given by applying $\text{pr}\hat{X}$ to Eq. (1.1) and then eliminating u_t , \bar{u}_t , and x -derivatives of u_t , \bar{u}_t through Eq. (1.1), its complex conjugate equation, and differential consequences of these equations. Here D_t

and D_x denote total derivatives with respect to t and x . Thus, a solution $u = f(t, x)$ of the Schrödinger equation (1.1) is *group-invariant* under a one-dimensional point symmetry group with a generator (2.1) if (and only if) it satisfies the additional equation

$$\eta(t, x, u, \bar{u}) - \tau(t, x, u, \bar{u})u_t - \zeta(t, x, u, \bar{u}) \cdot \nabla u = 0. \quad (2.5)$$

It is straightforward to solve Eq. (2.4) to determine all point symmetry generators (2.1). In particular, because τ, ζ, η do not depend on any derivatives of u and \bar{u} , Eq. (2.4) splits with respect to x -derivatives of u and \bar{u} , yielding an overdetermined linear system of PDEs on the functions $\tau(t, x, u, \bar{u}), \zeta(t, x, u, \bar{u}), \eta(t, x, u, \bar{u})$. This system leads to the following well-known result.^{6,10} (Here \odot and \wedge will, respectively, denote the symmetric and antisymmetric parts of an outer product.)

Theorem 1. *The point symmetries of the Schrödinger equation (1.1) are generated by*

$$\text{phase rotation} \quad X_1 = iu\partial/\partial u - i\bar{u}\partial/\partial \bar{u}, \quad (2.6a)$$

$$\text{time translation} \quad X_2 = \partial/\partial t, \quad (2.6b)$$

$$\text{space translations} \quad X_{3(l)} = e_{(l)} \cdot \partial/\partial x, \quad l = 1, \dots, n, \quad (2.6c)$$

$$\text{Galilean boosts} \quad X_{4(l)} = 2te_{(l)} \cdot \partial/\partial x - i(e_{(l)} \cdot x)(u\partial/\partial u - \bar{u}\partial/\partial \bar{u}), \quad (2.6d)$$

$$l = 1, \dots, n,$$

$$\text{space rotations} \quad X_{5(l,l')} = (e_{(l)} \cdot x)e_{(l')} \cdot \partial/\partial x - (e_{(l')} \cdot x)e_{(l)} \cdot \partial/\partial x, \quad (2.6e)$$

$$l = 1, \dots, n-1; l' = l+1, \dots, n,$$

$$\text{scaling} \quad X_6 = 2t\partial/\partial t + x \cdot \partial/\partial x - (2/p)u\partial/\partial u - (2/p)\bar{u}\partial/\partial \bar{u}, \quad (2.6f)$$

$$\begin{aligned} \text{inversion} \quad X_7 = t^2\partial/\partial t + tx \cdot \partial/\partial x - (2t/p + i|x|^2/4)u\partial/\partial u \\ - (2t/p - i|x|^2/4)\bar{u}\partial/\partial \bar{u} \quad \text{only for } p = 4/n, \end{aligned} \quad (2.6g)$$

where $\{e_{(1)}, \dots, e_{(n)}\}$ is any orthonormal basis for \mathbb{R}^n . The corresponding transformation groups acting on solutions $u = f(t, x)$ of the Schrödinger equation (1.1) are given by

$$u = \exp(i\phi)f(t, x), \quad (2.7a)$$

$$u = f(t - \epsilon, x), \quad (2.7b)$$

$$u = f(t, x - \epsilon e_{(l)}), \quad l = 1, \dots, n, \quad (2.7c)$$

$$u = \exp(-i(2\epsilon e_{(l)} \cdot x - \epsilon^2 t)/4)f(t, x - \epsilon t e_{(l)}), \quad l = 1, \dots, n, \quad (2.7d)$$

$$u = f(t, x + (\cos(\phi) - 1)(e_{(l)} \odot e_{(l)} + e_{(l')} \odot e_{(l')}) \cdot x + 2\sin(\phi)(e_{(l)} \wedge e_{(l')}) \cdot x), \quad (2.7e)$$

$$l = 1, \dots, n-1; l' = l+1, \dots, n,$$

$$u = \lambda^{-2/p}f(\lambda^{-2}t, \lambda^{-1}x), \quad (2.7f)$$

$$u = (1 + \epsilon t)^{-2/p} \exp(-i\epsilon|x|^2/(4 + 4\epsilon t))f(t/(1 + \epsilon t), x/(1 + \epsilon t)) \quad (2.7g)$$

$$\text{only for } p = 4/n,$$

with group parameters $-\infty < \epsilon < \infty, 0 < \lambda < \infty, 0 \leq \phi < 2\pi$.

The special power $p = 4/n$ for which the inversion group exists is commonly called the *pseudo-conformal power*. In the case $p \neq 4/n$, the transformations (2.7a)–(2.7f) comprise a semi-direct product of a scaling group acting on a central extension of the Galilean group, which has the Lie algebra structure

$$[X_1, X_2] = [X_1, X_{3(l)}] = [X_1, X_{4(l)}] = [X_1, X_{5(l,l')}] = [X_1, X_6] = 0, \quad (2.8)$$

$$[X_2, X_{3(l)}] = 0, \quad [X_2, X_{4(l)}] = 2X_{3(l)}, \quad [X_2, X_{5(l,l')}] = 0, \quad (2.9)$$

$$[X_{3(k)}, X_{4(l)}] = -\delta_{kl}X_1, \quad [X_{3(k)}, X_{5(l,l')}] = \delta_{kl}X_{3(l')} - \delta_{kl'}X_{3(l)}, \quad (2.10)$$

$$[X_{4(k)}, X_{5(l,l')}] = \delta_{kl}X_{4(l')} - \delta_{kl'}X_{4(l)}, \quad (2.11)$$

$$[X_6, X_2] = -2X_2, \quad [X_6, X_{3(k)}] = -X_{3(k)}, \quad [X_6, X_{4(k)}] = 2X_{4(k)}, \quad [X_6, X_{5(l,l')}] = 0. \quad (2.12)$$

In the case $p = 4/n$, the inversion transformation (2.7g) intertwines non-trivially with the previous group, as given by the commutator structure

$$[X_7, X_1] = [X_7, X_{4(l)}] = [X_7, X_{5(l,l')}] = 0, \quad (2.13)$$

$$[X_7, X_2] = -X_6, \quad [X_7, X_{3(l)}] = -\frac{1}{2}X_{4(l)}, \quad [X_7, X_6] = -2X_7. \quad (2.14)$$

For any domain $\Omega \subseteq \mathbb{R}^n$, the Schrödinger equation (1.1) has the variational formulation

$$\frac{\delta \mathcal{L}}{\delta \bar{u}} = 0 \quad (2.15)$$

given by the Lagrangian functional

$$\mathcal{L} = \int_{t_0}^{t_1} \int_{\Omega} L(u, \bar{u}, u_t, \bar{u}_t, \nabla u, \nabla \bar{u}) d^n x dt, \quad L = \frac{1}{2}i\bar{u}u_t - \frac{1}{2}iu\bar{u}_t + |\nabla u|^2 - \frac{2}{p+2}k|u|^{2+p}. \quad (2.16)$$

A *variational point symmetry* of this functional (2.16) is an infinitesimal point transformation (2.1) on (t, x, u, \bar{u}) under which \mathcal{L} is invariant up to spatial boundary terms at $\partial\Omega$ and temporal boundary terms at $t = t_0$ and $t = t_1$. This invariance condition holds if and only if the Lagrangian satisfies

$$\text{pr}\hat{X}L = D_t A + D_x \cdot B \quad (2.17)$$

for some functions A and B of t, x, u, \bar{u} , and derivatives of u, \bar{u} with respect to t and x , where \hat{X} is the characteristic form (2.2) and (2.3) of the generator (2.1). An equivalent condition on the Lagrangian is that $\text{pr}\hat{X}L$ is annihilated by the variational derivatives with respect to u and \bar{u} .

Since invariance of \mathcal{L} implies that its extrema (2.15) are preserved, every variational point symmetry of the Lagrangian functional (2.16) for the Schrödinger equation (1.1) is thereby a point symmetry of the Schrödinger equation itself such that

$$\frac{\delta(\text{pr}\hat{X}L)}{\delta \bar{u}} = 0, \quad \frac{\delta(\text{pr}\hat{X}L)}{\delta u} = 0. \quad (2.18)$$

This provides a straightforward way to determine all of the variational point symmetries starting from Theorem 1, which yields the following result.

Theorem 2. *The variational point symmetries of the Schrödinger equation (1.1) are generated by the symmetries (2.6a)–(2.6e) for all powers p and additionally by the symmetries (2.6f) and (2.6g) for the pseudo-conformal power $p = 4/n$.*

Thus, among all of the point symmetries listed in Theorem 1, the only non-variational symmetry is the scaling (2.6f) for $p \neq 4/n$.

Variational symmetries give rise to conservation laws for the Schrödinger equation (1.1) by means of Noether's theorem as follows. The invariance condition (2.17) combined with the variational identity

$$\text{pr}\hat{X}L = \frac{\delta L}{\delta u}P + \frac{\delta L}{\delta \bar{u}}\bar{P} + D_t\left(\frac{\partial L}{\partial u_t}P + \frac{\partial L}{\partial \bar{u}_t}\bar{P}\right) + D_x \cdot \left(\frac{\partial L}{\partial \nabla u}P + \frac{\partial L}{\partial \nabla \bar{u}}\bar{P}\right) \quad (2.19)$$

yields the multiplier equation

$$\text{Re}((iu_t - \Delta u - k|u|^p u)\bar{P}) = D_t T + D_x \cdot X \quad (2.20)$$

with

$$T = A - \frac{\partial L}{\partial u_t}P - \frac{\partial L}{\partial \bar{u}_t}\bar{P}, \quad X = B - \frac{\partial L}{\partial \nabla u}P - \frac{\partial L}{\partial \nabla \bar{u}}\bar{P}. \quad (2.21)$$

On all formal solutions $u(t, x)$ of the Schrödinger equation (1.1), the multiplier equation (2.20) then produces a *conservation law*

$$D_t T + D_x \cdot X = 0, \quad (2.22)$$

where the *conserved density* T and *flux* X are given by the Noether relation (2.21). In particular, A and B can be shown to have the specific form

$$A = -\tau L, \quad B = -\zeta L \quad (2.23)$$

as derived from the invariance condition (2.17). Therefore, the conserved density and flux are given explicitly by the simple formulas

$$T = -\tau L + \text{Re}(iu\bar{P}), \quad X = -\zeta L - 2\text{Re}(\bar{P}\nabla u) \quad (2.24)$$

in terms of L and $P = \eta - \tau u_t - \zeta \cdot \nabla u$. This result together with Theorem 2 yields all of the conservation laws generated by the variational point symmetries of the Schrödinger equation (1.1), as shown in Table I using the notation $\nabla_{(l)} = e_{(l)} \cdot \nabla$ where $\{e_{(1)}, \dots, e_{(n)}\}$ is any orthonormal basis for \mathbb{R}^n .

A. Radial reduction

A space rotation generator (2.6e) acts as an infinitesimal rotation in the 2-plane specified by a pair of basis vectors $e_{(l)}$ and $e_{(l')}$ in \mathbb{R}^n with $l \neq l'$. The corresponding group of transformations

$$(t, x, u) \rightarrow (t, x + (\cos(\phi) - 1)(e_{(l)} \odot e_{(l)} + e_{(l')} \odot e_{(l')}) \cdot x + 2 \sin(\phi)(e_{(l)} \wedge e_{(l')}) \cdot x, u) \quad (2.25)$$

is an $\text{SO}(2)$ Lie group. Composition of all such transformations acting in the $n(n-1)/2$ distinct 2-planes determined by an orthonormal basis $\{e_{(1)}, \dots, e_{(n)}\}$ of \mathbb{R}^n produces an $\text{SO}(n)$ Lie group of rotations, whose invariants consist of functions of t , $|x| = r$, u , \bar{u} .

Reduction of the Schrödinger equation (1.1) in \mathbb{R}^n under this point symmetry group of rotations gives the radial Schrödinger equation (1.2) where $u(t, r)$ is a group-invariant solution. The radial Lagrangian (1.10) arises naturally from this reduction of the Lagrangian functional (2.16) in \mathbb{R}^n , as given by

$$\mathcal{L}_{\text{rad.}} = \int_{t_0}^{t_1} \int_0^\infty L(r, u, \bar{u}, u_t, \bar{u}_t, u_r, \bar{u}_r) dr dt \quad (2.26)$$

in the radial domain $0 \leq r < \infty$.

From general results on symmetry reductions,⁴ a point symmetry generator (2.1) of the Schrödinger equation (1.1) in \mathbb{R}^n will admit a radial reduction if and only if it belongs to the normalizer algebra of the $\mathfrak{so}(n)$ Lie subalgebra of space rotations (2.6e) in the Lie algebra of all point symmetry generators (2.6). Likewise, through Noether's theorem, a conservation law (2.22) of the Schrödinger equation (1.1) will admit a radial reduction if and only if its corresponding variational point symmetry generator belongs to this same normalizer algebra.

Because the set of all formal radial solutions $u(t, r)$ is contained strictly as a subset in the set of all formal solutions $u(t, x)$ to the Schrödinger equation (1.1) in \mathbb{R}^n , the radial Schrödinger equation (1.2)

TABLE I. Conservation laws of the semilinear Schrödinger equation (1.1) in \mathbb{R}^n .

	T	X	
1	$- u ^2$	$i(u\nabla\bar{u} - \bar{u}\nabla u)$	L^2 norm
2	$ \nabla u ^2 - \frac{2}{p+2}k u ^{2+p}$	$-u_t\nabla\bar{u} - \bar{u}_t\nabla u$	Energy
3	$\frac{1}{2}i(u\nabla_{(l)}\bar{u} - \bar{u}\nabla_{(l)}u)$	$(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})e_{(l)}$ $+ \frac{1}{2}i(\bar{u}u_t - u\bar{u}_t)e_{(l)} - \nabla_{(l)}u\nabla\bar{u} - \nabla_{(l)}\bar{u}\nabla u$	Momentum $l = 1, \dots, n$
4	$it(u\nabla_{(l)}\bar{u} - \bar{u}\nabla_{(l)}u)$ $+ (e_{(l)} \cdot x) u ^2$	$-2t((\nabla\bar{u})(\nabla_{(l)}u) + (\nabla u)(\nabla_{(l)}\bar{u}))$ $+ 2t(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})e_{(l)}$ $+ i(\bar{u}u_t - u\bar{u}_t)e_{(l)} - i(e_{(l)} \cdot x)(u\nabla\bar{u} - \bar{u}\nabla u)$	Galilean momentum $l = 1, \dots, n$
5	$\frac{1}{2}i(e_{(l')} \cdot x)(\bar{u}\nabla_{(l)}u - u\nabla_{(l)}\bar{u})$ $- \frac{1}{2}i(e_{(l)} \cdot x)(\bar{u}\nabla_{(l')}u - u\nabla_{(l')} \bar{u})$	$(e_{(l')} \cdot x)(\nabla_{(l)}u\nabla\bar{u} + \nabla_{(l)}\bar{u}\nabla u)$ $-(e_{(l)} \cdot x)(\nabla_{(l')}u\nabla\bar{u} + \nabla_{(l')} \bar{u}\nabla u)$ $+ i(e_{(l)} \wedge e_{(l')}) \cdot x(\bar{u}u_t - u\bar{u}_t)$ $+ 2(e_{(l)} \wedge e_{(l')}) \cdot x(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})$	Angular momentum $l, l' = 1, \dots, n$ $l \neq l'$
6	$2t(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}i((x \cdot \nabla\bar{u})u - (x \cdot \nabla u)\bar{u})$	$(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})x$ $+ \frac{1}{2}i(\bar{u}u_t - u\bar{u}_t)x - \frac{2}{p}\nabla(u ^2)$ $- 2t(u_t\nabla\bar{u} + \bar{u}_t\nabla u) - (x \cdot \nabla u)\nabla\bar{u} - (x \cdot \nabla\bar{u})\nabla u$	Dilation energy $p = 4/n$
7	$t^2(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}it((x \cdot \nabla\bar{u})u - (x \cdot \nabla u)\bar{u})$ $+ \frac{1}{4} x ^2 u ^2$	$t(\nabla u ^2 - \frac{2}{p+2}k u ^{2+p})x$ $- \frac{1}{2}it(\bar{u}u_t - u\bar{u}_t)x - \frac{2}{p}t\nabla(u ^2)$ $- 2t((x \cdot \nabla u)\nabla\bar{u} + (x \cdot \nabla\bar{u})\nabla u)$ $- t^2(u_t\nabla\bar{u} + \bar{u}_t\nabla u) + \frac{1}{4} x ^2(\bar{u}\nabla u - u\nabla\bar{u})$	Pseudo-conformal energy $p = 4/n$

could possibly admit additional (“hidden”) point symmetries and conservation laws other than those inherited through radial reduction. Consequently, the only way to find all radial point symmetries, as well as all radial conservation laws corresponding to variational point symmetries, is by directly solving the determining equations for their generators.

For the radial Schrödinger equation (1.2), the generator of a point symmetry acting on the variables (t, r, u, \bar{u}) is given by

$$X = \tau(t, r, u, \bar{u})\partial/\partial t + \rho(t, r, u, \bar{u})\partial/\partial r + \eta(t, r, u, \bar{u})\partial/\partial u + \bar{\eta}(t, r, u, \bar{u})\partial/\partial \bar{u} \quad (2.27)$$

such that $\text{pr}X(iu_t - u_{rr} - (n-1)r^{-1}u_r - k|u|^p u) = 0$ holds for all formal solutions $u(t, r)$ of equation (1.2). The characteristic form for each such generator

$$\hat{X} = P(t, r, u, \bar{u}, u_t, \bar{u}_t, u_r, \bar{u}_r)\partial/\partial u + \bar{P}(t, r, u, \bar{u}, u_t, \bar{u}_t, u_r, \bar{u}_r)\partial/\partial \bar{u} \quad (2.28)$$

with

$$P = \eta - \tau u_t - \rho u_r \quad (2.29)$$

satisfies

$$iD_t P - D_r^2 P - (n-1)r^{-1}D_r P - k(1+p/2)\bar{u}^{p/2}u^{p/2}P - k(p/2)\bar{u}^{p/2-1}u^{1+p/2}\bar{P} = 0, \quad (2.30)$$

where u_t , \bar{u}_t , and r -derivatives of u_t , \bar{u}_t are eliminated through Eq. (1.2) and the complex-conjugate equation. Here D_t and D_r denote total derivatives with respect to t and r . This determining equation (2.30) splits with respect to r -derivatives of u and \bar{u} , giving an overdetermined linear system of PDEs

on the functions $\tau(t, r, u, \bar{u})$, $\rho(t, r, u, \bar{u})$, $\eta(t, r, u, \bar{u})$. After a straightforward integrability analysis, the system reduces to the PDEs

$$\tau_r = \tau_u = \tau_{\bar{u}} = \rho_u = \rho_{\bar{u}} = \eta_{\bar{u}} = \bar{\eta}_u = 0, \quad (2.31a)$$

$$\eta_{uu} = \bar{\eta}_{\bar{u}\bar{u}} = (pn - 4)\tau_{tt} = \tau_{ttt} = 0, \quad (2.31b)$$

$$\rho = \frac{1}{2}r\tau_t, \quad \eta_{ut} = \bar{\eta}_{\bar{u}t}, \quad \eta_{ur} - \bar{\eta}_{\bar{u}r} = -i\rho_t, \quad \eta_u + \bar{\eta}_{\bar{u}} = -\frac{4}{p}\rho_r, \quad (2.31c)$$

$$\bar{u}\eta + u\bar{\eta} = \bar{u}u(\eta_u + \bar{\eta}_{\bar{u}}), \quad (2.31d)$$

which are easily solved to obtain the following result.

Theorem 3. *The point symmetries of the radial Schrödinger equation (1.2) are generated by*

$$\text{phase rotation} \quad X_{\text{phas.}} = iu\partial/\partial u - i\bar{u}\partial/\partial \bar{u}, \quad (2.32a)$$

$$\text{time translation} \quad X_{\text{trans.}} = \partial/\partial t, \quad (2.32b)$$

$$\text{scaling} \quad X_{\text{scal.}} = 2t\partial/\partial t + r\partial/\partial r - (2/p)u\partial/\partial u - (2/p)\bar{u}\partial/\partial \bar{u}, \quad (2.32c)$$

$$\begin{aligned} \text{inversion} \quad X_{\text{inver.}} = & t^2\partial/\partial t + tr\partial/\partial r - (2t/p + ir^2/4)u\partial/\partial u \\ & - (2t/p - ir^2/4)\bar{u}\partial/\partial \bar{u} \quad \text{only for } p = 4/n, \end{aligned} \quad (2.32d)$$

with the Lie algebra structure

$$[X_{\text{phas.}}, X_{\text{trans.}}] = [X_{\text{phas.}}, X_{\text{scal.}}] = [X_{\text{phas.}}, X_{\text{inver.}}] = 0, \quad (2.33)$$

$$[X_{\text{trans.}}, X_{\text{scal.}}] = 2X_{\text{trans.}}, \quad [X_{\text{trans.}}, X_{\text{inver.}}] = X_{\text{scal.}}, \quad [X_{\text{scal.}}, X_{\text{inver.}}] = 2X_{\text{inver.}}. \quad (2.34)$$

The corresponding transformation groups acting on solutions $u = f(t, r)$ are given by

$$u = \exp(i\phi)f(t, r), \quad (2.35a)$$

$$u = f(t - \epsilon, r), \quad (2.35b)$$

$$u = \lambda^{-2/p}f(\lambda^{-2}t, \lambda^{-1}r), \quad (2.35c)$$

$$u = (1 + \epsilon t)^{-2/p} \exp(-i\epsilon r^2/(4 + 4\epsilon t))f(t/(1 + \epsilon t), r/(1 + \epsilon t)) \quad \text{only for } p = 4/n, \quad (2.35d)$$

with group parameters $-\infty < \epsilon < \infty$, $0 < \lambda < \infty$, $0 \leq \phi < 2\pi$.

Thus, the only point symmetries admitted by the radial Schrödinger equation (1.2) are the ones it inherits from reduction of the Schrödinger equation (1.2) in \mathbb{R}^n under the $SO(n)$ group of rotations.

The Lagrangian functional (2.26) for the radial Schrödinger equation (1.2) is invariant up to spatial boundary terms (at $r = 0$ and $r \rightarrow \infty$) and temporal boundary terms (at $t = t_0$ and $t = t_1$) under an infinitesimal point transformation (2.27) on (t, r, u, \bar{u}) if and only if

$$\text{pr}\hat{X}L = D_t A + D_r B \quad (2.36)$$

holds for some functions A and B of t, r, u, \bar{u} , and derivatives of u, \bar{u} with respect to t and r . This is the condition for X to be a variational point symmetry of the Lagrangian functional (2.26). Since the transformation group generated by a variational point symmetry necessarily preserves the extrema $\delta\mathcal{L}_{\text{rad.}}/\delta\bar{u} = 0$ of this functional $\mathcal{L}_{\text{rad.}}$, every such symmetry generator X is thereby a point

TABLE II. Conservation laws of the semilinear radial Schrödinger equation (1.2).

	T	X	
1	$-r^{n-1} u ^2$	$ir^{n-1}(u\bar{u}_r - \bar{u}u_r)$	L^2 norm (charge)
2	$r^{n-1}(u_r ^2 - \frac{2}{p+2}k u ^{2+p})$	$-r^{n-1}(u_t\bar{u}_r + \bar{u}_t u_r)$	Energy
3	$2tr^{n-1}(u_r ^2 - \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}ir^n(u\bar{u}_r - \bar{u}u_r)$	$-r^n(u_r ^2 + \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}ir^n(\bar{u}u_t - u\bar{u}_t) - \frac{2}{p}r^{n-1}(u\bar{u}_r + \bar{u}u_r)$ $- 2tr^{n-1}(u_t\bar{u}_r + \bar{u}_t u_r)$	Dilation energy $p = 4/n$
4	$t^2r^{n-1}(u_r ^2 - \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}itr^n(u\bar{u}_r - \bar{u}u_r)$ $+ \frac{1}{4}r^{n+1} u ^2$	$-\frac{2}{p}tr^{n-1}(u\bar{u}_r + \bar{u}u_r)$ $-t^2r^{n-1}(u_t\bar{u}_r + \bar{u}_t u_r)$ $-tr^n(u_r ^2 + \frac{2}{p+2}k u ^{2+p})$ $+ \frac{1}{2}itr^n(\bar{u}u_t - u\bar{u}_t) + \frac{1}{4}ir^{n+1}(\bar{u}u_r - u\bar{u}_r)$	Pseudo-conformal energy $p = 4/n$

symmetry of the radial Schrödinger equation itself for which the variational condition (2.36) holds, or equivalently for which the equation

$$\frac{\delta(\text{pr}\hat{X}L)}{\delta\bar{u}} = \frac{\delta(\text{pr}\hat{X}L)}{\delta u} = 0 \quad (2.37)$$

is satisfied. Condition (2.37) can be shown to reduce to the symmetry determining equations (2.31) plus the additional equation

$$\tau_t + (n-1)r^{-1}\rho - (1+4/p)\rho_r = 0. \quad (2.38)$$

It is straightforward to determine which of the point symmetries from Theorem 3 satisfy Eq. (2.38), which yields the following result.

Theorem 4. *The variational point symmetries of the radial Schrödinger equation (1.2) are generated by the symmetries (2.32a) and (2.32b) when $p \neq 4/n$ and by all of the symmetries (2.32a)–(2.32d) when $p = 4/n$.*

From Noether's theorem, these variational symmetries yield the multiplier equation

$$\text{Re}((iu_t - u_{rr} - (n-1)r^{-1}u_r - k|u|^p u)\bar{P}) = D_t T + D_r X, \quad (2.39)$$

which produces a conservation law

$$D_t T + D_r X = 0 \quad (2.40)$$

holding on all formal solutions $u(t, r)$ of the radial Schrödinger equation (1.2). Here P is the function (2.29) given by the characteristic form of a symmetry generator (2.28), while the conserved density T and flux X are given by the simple formulas

$$T = -\tau L + \text{Re}(iu\bar{P}), \quad X = -\rho L - 2\text{Re}(\bar{P}u_r), \quad (2.41)$$

which can be derived from the variational condition (2.36) in terms of the radial Lagrangian (1.10). The resulting conservation laws generated by all of the variational point symmetries in Theorem 4 are shown in Table II.

B. Group invariance of radial solutions

A solution $u = f(t, r)$ of the radial Schrödinger equation (1.2) is *group-invariant* under a one-dimensional point symmetry group if (and only if) it satisfies $\hat{X}u|_{u=f(t,r)} = 0$ where \hat{X} is the symmetry generator in characteristic form (2.28) and (2.29). To find all group-invariant solutions, it is sufficient to consider a maximal set of one-dimensional point symmetry subgroups that are

conjugacy inequivalent in the full Lie group of point symmetries admitted by the radial Schrödinger equation (1.2). This is most easily carried out at the Lie algebra level by exhibiting a maximal set of one-dimensional subalgebras that are conjugacy inequivalent in the Lie algebra of point symmetry generators. Such a set, called *optimal symmetry generators*, can be straightforwardly determined by the methods in Refs. 4 and 11 starting from the point symmetry algebra given in Theorem 3.

Lemma 1. The Lie algebra of point symmetry generators (2.32) for the radial Schrödinger equation (1.2) is isomorphic to $\mathfrak{u}(1) \oplus A(2)$ when $p \neq 4/n$ and $\mathfrak{u}(1) \oplus \mathfrak{sl}(2, \mathbb{R})$ when $p = 4/n$. (Here $A(2)$ denotes a 2-dimensional, non-abelian Lie algebra, which is unique up to isomorphism.) An optimal set of symmetry generators X consists of

$$X_{\text{phas.}}, \quad (2.42a)$$

$$X_{\text{trans.}} + \nu X_{\text{phas.}}, \quad (2.42b)$$

$$X_{\text{scal.}} + \mu X_{\text{phas.}}, \quad (2.42c)$$

with parameters $-\infty < \nu < \infty$, $-\infty < \mu < \infty$, and also

$$X_{\text{trans.}} + X_{\text{inver.}} + \kappa X_{\text{phas.}} \quad \text{for } p = 4/n \quad (2.42d)$$

with parameter $-\infty < \kappa < \infty$.

These generators (2.42) define group-invariant solutions whose form is determined by the respective characteristic equations

$$u = 0, \quad (2.43a)$$

$$i\nu u - u_t = 0, \quad (2.43b)$$

$$(\frac{2}{p} - i\mu)u + 2tu_t + ru_r = 0, \quad (2.43c)$$

$$(\frac{2}{p}t + i(\frac{1}{4}r^2 - \kappa))u + (t^2 + 1)u_t + tru_r = 0, \quad p = 4/n, \quad (2.43d)$$

which come from $\hat{X}u = 0$, where \hat{X} denotes a generator (2.42) in characteristic form (2.28) and (2.29). The corresponding one-dimensional point transformation subgroups are given by

$$t \rightarrow t, r \rightarrow r, u \rightarrow \exp(i\phi)u, \quad (2.44a)$$

$$t \rightarrow t + \epsilon, r \rightarrow r, u \rightarrow \exp(i\nu\epsilon)u, \quad (2.44b)$$

$$t \rightarrow \lambda^2 t, r \rightarrow \lambda r, u \rightarrow \exp(i\mu \ln \lambda)\lambda^{-2/p}u, \quad (2.44c)$$

$$t \rightarrow (\sin \phi + t \cos \phi)/(\cos \phi - t \sin \phi), r \rightarrow r/(\cos \phi - t \sin \phi), \\ u \rightarrow \exp(i\kappa\phi - i\frac{1}{4}r^2 \sin \phi/(\cos \phi - t \sin \phi))(\cos \phi - t \sin \phi)^{2/p}u, \quad p = 4/n, \quad (2.44d)$$

with group parameters $-\infty < \epsilon < \infty$, $0 < \lambda < \infty$, $0 \leq \phi < 2\pi$.

Invariance under the phase rotation symmetry generator (2.42a) yields only a trivial group-invariant solution (2.43a). Hence, for finding non-trivial group-invariant solutions, only the remaining symmetry generators (2.42b)–(2.42d) in the optimal set need to be considered. The specific form for these solutions is given by integration of the characteristic equations (2.43b)–(2.43d), which can be expressed directly in terms of the invariants $\xi = \Xi(t, r)$, $U = \Upsilon(t, r, u)$, $\bar{U} = \bar{\Upsilon}(t, r, \bar{u})$ of the corresponding generators X .

C. Radial scaling reductions

The symmetry group of scaling transformations (2.44c) with the generator (2.42c) has invariants

$$\xi = t/r^2, \quad U = r^{2/p} \exp(-i\mu \ln r)u, \quad \bar{U} = r^{2/p} \exp(i\mu \ln r)\bar{u}, \quad (2.45)$$

depending on a parameter $-\infty < \mu < \infty$. Hence the corresponding form for group-invariant solutions (2.43c) is given by

$$u(t, r) = r^{-2/p} \exp(i\mu \ln r)U(\xi). \quad (2.46)$$

The radial Schrödinger equation (1.2) thereby reduces to the complex semilinear 2nd order ODE

$$4\xi^2 U'' + ((8 - 2n + 8/p)\xi - i(1 + 4\mu\xi))U' + ((4 - 2n)/p + 4/p^2 - \mu^2 + i\mu(n - 2 - 4/p))U + k|U|^p U = 0. \quad (2.47)$$

For $\mu = 0$, this reduction describes radial similarity solutions.

From Theorem 4, the symmetry group (2.44c) leaves invariant the Lagrangian functional (2.26) for the radial Schrödinger equation if and only if $p = 4/n$. Hence, in the case of the pseudo-conformal power, ODE (2.47) has a variational formulation given by reduction of the radial Schrödinger Lagrangian (1.10),

$$L = -4\xi^2 |U'|^2 + \left(\frac{4(p-1)}{p^2} - \mu^2\right)|U|^2 + \frac{2}{p+2}k|U|^{2+p} + i\left(\frac{1}{2} + 2\mu\xi\right)(U\bar{U}' - \bar{U}U') \text{ for } p = 4/n. \quad (2.48)$$

D. Radial time-translation reductions

The symmetry group of time translation transformations (2.44b) with the generator (2.42b) has invariants

$$\xi = r, \quad U = \exp(-ivt)u, \quad \bar{U} = \exp(ivt)\bar{u}, \quad (2.49)$$

depending on a parameter $-\infty < v < \infty$. Hence

$$u(t, r) = \exp(ivt)U(\xi) \quad (2.50)$$

yields the corresponding form for group-invariant solutions (2.43b). The radial Schrödinger equation (1.2) thereby reduces to the complex semilinear 2nd order ODE

$$U'' + (n - 1)\xi^{-1}U' + vU + k|U|^p U = 0. \quad (2.51)$$

For $v = 0$, this reduction describes radial static solutions, which are invariant under the symmetry subgroup of time translation transformations

$$t \rightarrow t + \epsilon, \quad r \rightarrow r, \quad u \rightarrow u. \quad (2.52)$$

When $v \neq 0$, the reduction (2.50) describes radial standing wave solutions.

Since, by Theorem 4, the Lagrangian functional (2.26) for the radial Schrödinger equation is invariant under the symmetry group (2.44b), ODE (2.51) has a variational formulation given by reduction of the radial Schrödinger Lagrangian (1.10),

$$L = \xi^{n-1}(-|U'|^2 + v|U|^2 + \frac{2}{p+2}k|U|^{2+p}). \quad (2.53)$$

E. Radial pseudo-conformal reductions

The symmetry group of combined inversion and time translation transformations (2.44d) with the generator (2.42d) for $p = 4/n$ has invariants

$$\begin{aligned} \xi &= (1 + t^2)/r^2, \quad U = r^{2/p} \exp(i\kappa \arctan(1/t) + ir^2 t/(4(1 + t^2)))u, \\ \bar{U} &= r^{2/p} \exp(-i\kappa \arctan(1/t) - ir^2 t/(4(1 + t^2)))\bar{u}, \end{aligned} \quad (2.54)$$

depending on a parameter $-\infty < \kappa < \infty$. Hence

$$u(t, r) = r^{-2/p} \exp(-i\kappa \arctan(1/t) - ir^2 t / (4(1+t^2))) U(\xi) \quad (2.55)$$

yields the corresponding form for group-invariant solutions (2.43d). The radial Schrödinger equation (1.2) thereby reduces to the complex semilinear 2nd order ODE

$$4\xi^2 U'' + 8\xi U' + (\kappa \xi^{-1} - \xi^{-2}/4 + n(1 - n/4))U + k|U|^{4/n}U = 0. \quad (2.56)$$

From Theorem 4, since the Lagrangian functional (2.26) for the radial Schrödinger equation is invariant under the symmetry group (2.44d), ODE (2.56) has a variational formulation given by reduction of the radial Schrödinger Lagrangian (1.10) with $p = 4/n$,

$$L = -4\xi^2 |U'|^2 + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} + n(1 - \frac{1}{4}n))|U|^2 + \frac{n}{n+2}k|U|^{2+4/n}. \quad (2.57)$$

F. Group invariance of optimal radial reductions

The ODEs (2.47), (2.51), and (2.56) arising by the reduction of the radial Schrödinger equation under its optimal point symmetry subgroups (2.42c), (2.42b), and (2.42d) have the following symmetry structure.

Proposition 1. (i) The point symmetries admitted by the scaling-group ODE (2.47) and the pseudo-conformal-group ODE (2.56) are generated only by phase rotations

$$Y_{\text{phas.}} = iU \partial / \partial U - i\bar{U} \partial / \partial \bar{U}. \quad (2.58)$$

(ii) The point symmetries admitted by the translation-group ODE (2.51) are generated by

$$\text{scaling } Y_{\text{scal.}} = \xi \partial / \partial \xi - (2/p)U \partial / \partial U - (2/p)\bar{U} \partial / \partial \bar{U} \quad \text{for } v = 0, \quad (2.59)$$

$$\text{dilation } Y_{\text{dil.}} = \xi^{2-n}(\xi / (n-2) \partial / \partial \xi - U \partial / \partial U - \bar{U} \partial / \partial \bar{U}) \quad (2.60)$$

$$\text{for } v = 0, \quad p = 2(3-n)/(n-2) \neq 0,$$

in addition to phase rotations (2.58).

In the case $v = 0$, the translation-group ODE (2.51) determines all static solutions of the radial Schrödinger equation (1.2),

$$U'' + (n-1)\xi^{-1}U' + k|U|^p U = 0 \quad \text{with } U = u, \xi = r. \quad (2.61)$$

The scaling symmetry (2.59) of this ODE is inherited from the scaling invariance of the radial Schrödinger equation, due to the commutator structure $[X_{\text{trans.}}, X_{\text{scal.}}] = 2X_{\text{trans.}}$. In contrast, the dilation symmetry (2.60) is not inherited from any invariance of the radial Schrödinger equation and thus describes a hidden symmetry arising for static solutions.

III. QUADRATURE OF COMPLEX 2ND ORDER SEMILINEAR ODES

Details of the two reduction of order methods outlined in Sec. I for solving complex 2nd order semilinear ODEs of the general form (2.47), (2.51), and (2.56) will now be presented.

A. Reduction by point symmetries

Consider a complex 2nd order semilinear ODE

$$\alpha(\xi)U'' + \beta(\xi)U' + \gamma(\xi)U + k|U|^p U = 0 \quad (3.1)$$

with a real independent variable ξ and a complex dependent variable U . Here α, β, γ are allowed to be complex functions of ξ , while k and p are assumed to be non-zero real constants. Every such

ODE (3.1) is invariant under a U(1) group of phase rotation symmetries given by the generator (2.58). Now suppose an ODE (3.1) admits another one-dimensional group of point symmetries, with a generator of the form

$$Y = \zeta(\xi)\partial/\partial\xi + \Omega(\xi)U\partial/\partial U + \bar{\Omega}(\xi)\bar{U}\partial/\partial\bar{U}, \quad (3.2)$$

whereby

$$[Y_{\text{phas.}}, Y] = 0. \quad (3.3)$$

If the ODE has a variational formulation for which both $Y_{\text{phas.}}$ and Y are variational symmetries, then these two symmetries can be used to reduce the ODE to quadratures by means of first integrals. However, the standard reduction steps⁴ would require finding the Lagrangian functional and checking its invariance under the two symmetries, which can be cumbersome to carry out. A simpler, more direct way to accomplish the same reduction is to utilize canonical coordinates and integrating factors associated to the symmetries $Y_{\text{phas.}}$ and Y as follows.

Make a change of variables from (ξ, U, \bar{U}) to (z, V, \bar{V}) given by the canonical coordinates of the generator (3.2),

$$Yz = 1, \quad YV = 0, \quad Y\bar{V} = 0, \quad (3.4)$$

where

$$z = \int (1/\zeta)d\xi, \quad V/U = \exp\left(-\int (\Omega/\zeta)d\xi\right) \quad (3.5)$$

are functions only of ξ . In terms of these variables, the generator (3.2) takes the form of a z -translation

$$Y = \partial/\partial z, \quad (3.6)$$

while the phase rotation generator (2.58) is given by

$$Y_{\text{phas.}} = iV\partial/\partial V - i\bar{V}\partial/\partial\bar{V}. \quad (3.7)$$

Consequently, the ODE (3.1) is transformed into

$$aV'' + bV' + cV + k|V|^p V = 0 \quad (3.8)$$

for $V(z)$, where the coefficients a, b, c are constants. The following result is straightforward to prove by considering the characteristic form of the symmetry generators (3.6) and (3.7),

$$\hat{Y} = -V'\partial/\partial V - \bar{V}'\partial/\partial\bar{V}, \quad \hat{Y}_{\text{phas.}} = iV\partial/\partial V - i\bar{V}\partial/\partial\bar{V}. \quad (3.9)$$

Lemma 2. (i) An ODE (3.8) has a non-trivial invariant solution with respect to z -translations if and only if $\bar{c} = c$ and $c/k < 0$ (unless p is a rational number having an odd numerator). Then $V' = 0$ yields

$$|V| = (-c/k)^{1/p}, \quad (3.10)$$

which determines V up to an arbitrary constant phase.

(ii) An ODE (3.8) has integrating factors \bar{V}' and $-i\bar{V}$ if and only if $\bar{a} = a, \bar{b} = -b, \bar{c} = c$. Then the first integrals are, respectively, given by

$$a|V'|^2 + c|V|^2 + 2k \int |V|^{p+1}d|V| = C_1, \quad (3.11)$$

$$a(\bar{V}V' - \bar{V}'V) + b|V|^2 = i2C_2, \quad (3.12)$$

from which $V(z)$ can be determined by quadratures, where C_1 and C_2 are arbitrary real constants.

The solutions arising from Lemma 2 can be easily obtained in a more explicit form in terms of polar variables

$$V = A \exp(i\Phi). \quad (3.13)$$

In particular, the first integrals (3.11) and (3.12) yield

$$aA'^2 = -aA^2\Phi'^2 - cA^2 - 2k \int A^{p+1} dA + C_1, \quad (3.14)$$

$$a\Phi' = \tilde{b} + C_2/A^2, \quad \tilde{b} = i2b. \quad (3.15)$$

Combining these two differential equations, we get the quadratures

$$\int \frac{dA}{\sqrt{H(A)}} = \pm z + C_3, \quad (3.16a)$$

$$\Phi = (\tilde{b}/a)z \pm (C_2/a) \int \frac{dA}{A^2\sqrt{H(A)}} + C_4, \quad (3.16b)$$

assuming $A' \neq 0$, where

$$H(A) = (aC_1 - 2\tilde{b}C_2)/a^2 - ((\tilde{b}/a)^2 + c/a)A^2 - (C_2/a)^2 A^{-2} - (2k/a) \int A^{p+1} dA. \quad (3.17)$$

The quadratures (3.16) determine the general solution $V(z)$ of ODE (3.8) in the case $\bar{a} = a$, $\bar{b} = -b$, $\bar{c} = c$. Moreover, these conditions on a , b , c are necessary and sufficient for the ODE (3.8) to have an Euler–Lagrange structure

$$\delta L/\delta \bar{V} = aV'' + bV' + cV + k|V|^p V = 0, \quad \bar{a} = a, \bar{b} = -b, \bar{c} = c \quad (3.18)$$

given by the Lagrangian

$$L = -a\bar{V}'V' + \frac{1}{2}b(\bar{V}V' - \bar{V}'V) + \frac{1}{2}c|V|^2 + 2k \int |V|^{p+1} d|V|, \quad (3.19)$$

for which both the z -translation generator (3.6) and the phase rotation generator (3.7) are variational symmetries. Noether's theorem thereby asserts

$$\text{pr}\hat{Y}L = -\frac{dL}{dz} = -V'\frac{\delta L}{\delta V} - \bar{V}'\frac{\delta L}{\delta \bar{V}} + \frac{d}{dz}\left((a\bar{V}' - \frac{1}{2}b\bar{V})V' + (aV' + \frac{1}{2}bV)\bar{V}'\right), \quad (3.20)$$

$$\text{pr}\hat{Y}_{\text{phas.}}L = 0 = iV\frac{\delta L}{\delta V} - i\bar{V}\frac{\delta L}{\delta \bar{V}} + \frac{d}{dz}\left((-a\bar{V}' + \frac{1}{2}b\bar{V})iV + (aV' + \frac{1}{2}bV)i\bar{V}\right), \quad (3.21)$$

giving a derivation of the first integrals (3.11) and (3.12).

In the general case without conditions on a , b , c , the z -translation invariant solution $V = \text{const}$ of ODE (3.8) is given by

$$A = (-c/k)^{1/p} = \text{const}, \quad \Phi = \text{const}. \quad (3.22)$$

Finally, changing variables (z, V, \bar{V}) back into (ξ, U, \bar{U}) in the original complex semilinear ODE (3.1) yields (i) the group-invariant solution with respect to the symmetry generator (3.2); (ii) the general solution under necessary and sufficient conditions for the generator (3.2) to be a variational symmetry.

B. Reduction by hidden conditional symmetries

Consider a real 1st-order Lagrangian of the form

$$L = \varpi(\xi)\left(-\tilde{\alpha}(\xi)\bar{U}'U' + i\tilde{\beta}(\xi)(\bar{U}'U - U'\bar{U}) + \tilde{\gamma}(\xi)|U|^2 + 2k \int |U|^{p+1} d|U|\right), \quad (3.23)$$

where ϖ , $\tilde{\alpha}$, $\tilde{\beta}$, $\tilde{\gamma}$ are real functions of ξ , and $k \neq 0$, $p \neq 0$ are real constants. The Euler–Lagrange equation $\delta L/\delta \bar{U} = 0$ yields a complex 2nd order semilinear ODE of the form (3.1) whose coefficients α , β , γ are given by

$$\alpha = \tilde{\alpha}, \quad \beta = \tilde{\alpha}' + \tilde{\alpha}\varpi'/\varpi - i2\tilde{\beta}, \quad \gamma = \tilde{\gamma} - i(\tilde{\beta}' + \tilde{\beta}\varpi'/\varpi). \quad (3.24)$$

Both the Lagrangian (3.23) and the Euler–Lagrange ODE (3.1) and (3.24) are invariant under the generator (2.58) of phase rotation symmetries on U, \tilde{U} . Hence $Y_{\text{phas.}}$ is a variational symmetry to which Noether's theorem can be applied to obtain a reduction of this ODE.

The simplest way to carry out the reduction is by use of polar variables

$$U = A \exp(i\Phi). \quad (3.25)$$

In particular, the ODE (3.1), (3.24) gets converted into a coupled semilinear system of real ODEs

$$\tilde{\alpha} A'' + ((\tilde{\alpha}\varpi)'/\varpi)A' - \tilde{\alpha}(\Phi' - \tilde{\beta}/\tilde{\alpha})^2 A + (\tilde{\gamma} + \tilde{\beta}^2/\tilde{\alpha})A + kA^{p+1} = 0, \quad (3.26)$$

$$\tilde{\alpha}\Phi'' A + 2\tilde{\alpha}(\Phi' - \tilde{\beta}/\tilde{\alpha})A' + ((\tilde{\alpha}\varpi)'/\varpi)A\Phi' - ((\tilde{\beta}\varpi)'/\varpi)A = 0, \quad (3.27)$$

which are seen to be the Euler–Lagrange equations $\delta L/\delta A = 0$ and $A^{-1}\delta L/\delta \Phi = 0$ of the Lagrangian (3.23) expressed in polar variables

$$L = \varpi(\xi) \left(-\tilde{\alpha}(\xi)A'^2 - \tilde{\alpha}(\xi)A^2(\Phi' - \tilde{\beta}(\xi)/\tilde{\alpha}(\xi))^2 + (\tilde{\gamma}(\xi) + \tilde{\beta}(\xi)^2/\tilde{\alpha}(\xi))A^2 + 2k \int A^{p+1} dA \right). \quad (3.28)$$

The phase rotation generator (2.58) thereby becomes

$$Y_{\text{phas.}} = \partial/\partial \Phi, \quad (3.29)$$

producing a group of shift transformations $\Phi \rightarrow \Phi + \epsilon$ with group parameter $-\infty < \epsilon < \infty$. Then the invariance $\text{pr}\hat{Y}_{\text{phas.}}L = 0$ yields the multiplier equation

$$\varpi A (\tilde{\alpha}\Phi'' A + 2\tilde{\alpha}(\Phi' - \tilde{\beta}/\tilde{\alpha})A' + ((\tilde{\alpha}\varpi)'/\varpi)\Phi' - ((\tilde{\beta}\varpi)'/\varpi)A) = \frac{d}{d\xi} \left(\varpi A^2 (\alpha\Phi' - \tilde{\beta}) \right), \quad (3.30)$$

which produces the first integral

$$\varpi A^2 (\tilde{\alpha}\Phi' - \tilde{\beta}) = C_1 \quad (3.31)$$

for the system (3.26) and (3.27). Through Eq. (3.31), Φ' can be eliminated in terms of A , so that the system reduces to a single real semilinear ODE

$$\tilde{\alpha} A'' + ((\tilde{\alpha}\varpi)'/\varpi)A' + (\tilde{\gamma} + \tilde{\beta}^2/\tilde{\alpha})A - (C_1^2/(\tilde{\alpha}\varpi^2))A^{-3} + kA^{p+1} = 0. \quad (3.32)$$

This ODE (3.32) is the Euler–Lagrange equation $\delta \tilde{L}/\delta A = 0$ of the modified Lagrangian

$$\begin{aligned} \tilde{L} &= \left(L + 2(\Phi' - \tilde{\beta}(\xi)/\tilde{\alpha}(\xi))C_1 \right) \Big|_{\tilde{\alpha}(\xi)\Phi' - \tilde{\beta}(\xi) = (C_1/\varpi(\xi))A^{-2}} \\ &= \varpi(\xi) \left(-\tilde{\alpha}(\xi)A'^2 + (\tilde{\gamma}(\xi) + \tilde{\beta}(\xi)^2/\tilde{\alpha}(\xi))A^2 + (C_1^2/(\tilde{\alpha}(\xi)\varpi(\xi)^2))A^{-2} + 2k \int A^{p+1} dA \right), \end{aligned} \quad (3.33)$$

where C_1 is treated as a constant parameter.

There are two cases, depending on C_1 and p , for which the ODE (3.32) can be solved. First, if $C_1 = 0$ and $p = -1$ then the ODE becomes linear

$$\tilde{\alpha} A'' + ((\tilde{\alpha}\varpi)'/\varpi)A' + (\tilde{\gamma} + \tilde{\beta}^2/\tilde{\alpha})A + k = 0 \quad (3.34)$$

and its general solution will be given by special functions. Second, if $C_1 \neq 0$ or $p \neq -1$ then the ODE will reduce to quadratures by means of integrating factors arising from variational symmetries of the Lagrangian (3.33).

Note that the set of solutions of the ODE (3.32) for each value of C_1 corresponds to the subset of solutions of the system (3.26) and (3.27) defined by the resulting level set of C_1 . Suppose a point symmetry generator

$$Y = \zeta(\xi)\partial/\partial \xi + \Upsilon(\xi)A\partial/\partial A + \omega(\xi)\partial/\partial \Phi \quad (3.35)$$

admitted by system (3.26) and (3.27) leaves invariant the level set of C_1 , so that $\text{pr}\hat{Y}C_1 = 2\varpi A(\Upsilon A - \zeta A')(\tilde{\alpha}\Phi' - \tilde{\beta}) + \varpi \tilde{\alpha} A^2(\omega - \zeta\Phi')' = 0$ when $\Phi' = (\tilde{\beta} + (C_1/\varpi)A^{-2})/\tilde{\alpha}$. Then, by

reduction, the level-set ODE (3.32) will inherit a corresponding point symmetry generator

$$\tilde{Y} = \zeta(\xi)\partial/\partial\xi + \Upsilon(\xi)A\partial/\partial A. \quad (3.36)$$

Such symmetries are equivalent to conditional point symmetries of the form (3.2) for the original complex 2nd order semilinear ODE (3.1) such that the generator preserves the first integral (3.31) given in terms of U, \bar{U} by

$$C_1 = \varpi(\tfrac{1}{2}i\tilde{\alpha}(U\bar{U}' - U'\bar{U}) - \tilde{\beta}U\bar{U}). \quad (3.37)$$

Any variational symmetries that are inherited in this way by the level-set ODE (3.32) will yield the same integrating factors derived from the previous reduction method in Sec. III A.

Now suppose that the level-set ODE (3.32) admits a variational point symmetry that is not inherited from any point symmetry admitted by the system (3.26) and (3.27). Such a variational symmetry will be a hidden conditional symmetry which yields an additional integrating factor for the ODE (3.32). The resulting reduction is most easily carried out by utilizing canonical coordinates associated to the symmetry generator as follows.

For a symmetry of the form (3.36), we change variables from (ξ, A) to canonical coordinates (z, B) given by

$$\tilde{Y}z = 1, \quad \tilde{Y}B = 0, \quad (3.38)$$

where

$$z = \int (1/\zeta)d\xi, \quad B = \exp\left(-\int (\Upsilon/\zeta)d\xi\right)A. \quad (3.39)$$

The symmetry generator (3.36) thereby becomes a z -translation

$$\tilde{Y} = \partial/\partial z \quad (3.40)$$

and thus the ODE (3.32) is transformed into

$$aB'' + bB' + cB + qB^{-3} + kB^{p+1} = 0 \quad (3.41)$$

for $B(z)$, where the coefficients a, b, c, q are real constants. Similarly, the Lagrangian (3.33) becomes (modulo a total z -derivative)

$$\zeta\tilde{L} = \hat{L} = \exp((b/a)z)\left(-aB'^2 + cB^2 - qB^{-2} + 2k\int B^{p+1}dB\right), \quad (3.42)$$

where $1/\zeta = \frac{dz}{d\xi}$ is the Jacobian of the transformation for the independent variable. Then Noether's theorem applied to the z -translation generator (3.40) in characteristic form

$$\widehat{\tilde{Y}} = -B'\partial/\partial B \quad (3.43)$$

yields the multiplier equation

$$\begin{aligned} \text{pr}\widehat{\tilde{Y}}(\hat{L}) &= -\exp((b/a)z)\frac{d}{dz}\left(-aB'^2 + cB^2 - qB^{-2} + 2k\int B^{p+1}dB\right) \\ &= -B'\frac{\delta\hat{L}}{\delta B} + \frac{d}{dz}\left(2(\exp((b/a)z)aB')B'\right). \end{aligned} \quad (3.44)$$

Equation (3.44) shows that $b = 0$ is a necessary and sufficient condition for the z -translation generator (3.40) to be a variational symmetry of the ODE (3.32). As an immediate consequence, the following result holds.

Lemma 3. An ODE (3.41) has integrating factor B' if and only if $b = 0$. The resulting first integral is given by

$$aB'^2 + cB^2 - qB^{-2} + 2k\int B^{p+1}dB = C_2, \quad (3.45)$$

from which $B(z)$ is determined by the quadrature

$$\int \frac{dB}{\sqrt{H(B)}} = \pm z + C_3, \quad (3.46)$$

where

$$H(B) = C_2/a - (c/a)B^2 + (q/a)B^{-2} - (2k/a) \int B^{p+1} dB. \quad (3.47)$$

When $b = 0$, the quadrature (3.46) combined with the change of variable (3.39) will yield the general solution of the ODE (3.32) for $A(\xi)$. In this case the first integral (3.31) provides a quadrature for $\Phi(\xi)$, thereby yielding the general solution of the system (3.26) and (3.27). These quadratures for $A(\xi)$ and $\Phi(\xi)$ can be written down in a fairly simple way. First, we find that the change of variable (3.39) relating the Lagrangians (3.42) and (3.33) gives

$$A^2/B^2 = \exp\left(2 \int (\Upsilon/\zeta) d\xi\right) = (a\zeta/(\varpi\tilde{\alpha})) \exp\left((b/a) \int (1/\zeta) d\xi\right). \quad (3.48)$$

Hence, from the quadrature (3.46) together with the variables (3.39), we see

$$\int \frac{dB}{\sqrt{H(B)}} = \pm \int (1/\zeta) d\xi + C_3 \quad (3.49a)$$

and

$$A = (a\zeta/(\varpi\tilde{\alpha}))^{1/2} B \quad (3.49b)$$

for $b = 0$. Last, changing variables in the first integral (3.31) and using the other first integral (3.45), we obtain

$$\Phi = \pm(C_1/a) \int \frac{dB}{B^2 \sqrt{H(B)}} + \int (\tilde{\beta}/\tilde{\alpha}) d\xi + C_4. \quad (3.49c)$$

The integrals (3.49) yield the general solution of the system (3.26) and (3.27) for $A(\xi)$ and $\Phi(\xi)$ under necessary and sufficient conditions for the reduced ODE (3.32) to admit a variational symmetry (3.36). This determines the corresponding solution $U(\xi) = A(\xi) \exp(i\Phi(\xi))$ to the original complex semilinear Euler–Lagrange ODE (3.1) and (3.24).

Finally, suppose the level-set ODE (3.32) admits a hidden point symmetry that is not a variational symmetry. Such a hidden conditional symmetry can be used to obtain an invariant solution of the ODE (3.32), which will determine a corresponding solution of the system (3.26) and (3.27) through the first integral (3.31) as follows.

Lemma 4. An ODE (3.32) has a non-trivial invariant solution with respect to z -translations if and only if $c \neq 0$ or $q \neq 0$ when $p \neq -4$, or $c \neq 0$ and $q \neq -k$ when $p = -4$. Then $B' = 0$ yields

$$0 = c + qB^{-4} + kB^p, \quad (3.50)$$

which is an algebraic equation determining $B = \text{const} \neq 0$.

Since $B' = 0$, relation (3.48) yields

$$A = (a\zeta/(\varpi\tilde{\alpha}))^{1/2} B \exp\left((b/2a) \int (1/\zeta) d\xi\right). \quad (3.51)$$

Then, after changing variables from (ξ, A) to (z, B) in the first integral (3.31), we find $\frac{d\Phi}{dz} = (C_1/(aB^2)) \exp(-(b/a)z) + \zeta \tilde{\beta}/\tilde{\alpha}$ from which we obtain the quadrature

$$\Phi = -C_1/(bB^2) \exp\left((-b/a) \int (1/\zeta) d\xi\right) + \int (\tilde{\beta}/\tilde{\alpha}) d\xi + C_2. \quad (3.52)$$

These expressions (3.51) and (3.52) yield a solution of the system (3.26) and (3.27), and hence $U(\xi) = A(\xi) \exp(i\Phi(\xi))$ gives a solution to the original Euler–Lagrange ODE (3.1) and (3.24),

corresponding to an invariant solution of the level-set ODE (3.32) under a hidden conditional point symmetry (3.36).

IV. SOLUTIONS TO THE OPTIMAL TRANSLATION-GROUP ODE

The reduction of order methods from Sec. III will now be applied to the translation-group ODE (2.51) arising by the reduction of the radial Schrödinger equation under its optimal subgroup of point symmetries (2.42b).

As a preliminary step, we use polar variables $U = A \exp(i\Phi)$ to convert this $U(1)$ -invariant ODE (2.51) into a semilinear system of real ODEs

$$A'' - A\Phi'^2 + (n-1)\xi^{-1}A' + \nu A + kA^{p+1} = 0, \quad (4.1a)$$

$$\Phi'' + 2A^{-1}A'\Phi' + (n-1)\xi^{-1}\Phi' = 0. \quad (4.1b)$$

The ODEs in this system (4.1) are the respective Euler–Lagrange equations $\delta L/\delta A = 0$ and $A^{-1}\delta L/\delta \Phi = 0$ of the $U(1)$ -invariant Lagrangian (2.53) expressed in polar variables,

$$L = \xi^{n-1}(-A'^2 - A^2\Phi'^2 + \nu A^2 + \frac{2}{2+p}kA^{2+p}). \quad (4.2)$$

In polar form, the point symmetries of ODE (2.51) listed in Proposition 1 consist of

$$\text{phase rotation } Y_{\text{phas.}} = \partial/\partial \Phi, \quad (4.3a)$$

$$\text{scaling } Y_{\text{scal.}} = \xi \partial/\partial \xi - (2/p)A \partial/\partial A \quad \text{for } \nu = 0, \quad (4.3b)$$

$$\text{dilation } Y_{\text{dil.}} = \xi^{2-n}(\xi \partial/\partial \xi + (2-n)A \partial/\partial A) \quad (4.3c)$$

for $\nu = 0, \quad p = 2(3-n)/(n-2).$

Invariance of the polar Lagrangian (4.2) under $Y_{\text{phas.}}$ produces a first integral (3.31), yielding

$$\Phi' = C_1 \xi^{1-n} A^{-2}. \quad (4.4)$$

The polar system (4.1) thereby reduces to a single real semilinear ODE

$$A'' + (n-1)\xi^{-1}A' + \nu A + kA^{1+p} - C_1^2 \xi^{2-2n} A^{-3} = 0, \quad (4.5)$$

which is the Euler–Lagrange equation $\delta \tilde{L}/\delta A = 0$ of a modified Lagrangian (3.33), given by

$$L = \xi^{n-1}(-A'^2 + \nu A^2 + \frac{2}{2+p}kA^{2+p} + C_1^2 \xi^{2-2n} A^{-2}). \quad (4.6)$$

Solutions of the ODE (4.5) for $A(\xi)$ represent the level set $C_1 = \text{const}$ of solutions $(A(\xi), \Phi(\xi))$ to the polar system (4.1), or equivalently the level set

$$C_1 = \frac{1}{2}i\xi^{n-1}(U\bar{U}' - U'\bar{U}) = \text{const} \quad (4.7)$$

of solutions $U(\xi)$ to the translation-group ODE (2.51).

Note that the level-set ODE (4.5) will be linear if and only if $C_1 = 0$ and $p = -1$.

Proposition 2. In the nonlinear case $C_1 \neq 0$ or $p \neq -1$, the level-set ODE (4.5) admits point symmetries only when

$$\nu = 0. \quad (4.8)$$

For this case the admitted point symmetries consist of

$$Y_1 = \xi \partial/\partial \xi - (2/p)A \partial/\partial A, \quad \text{for } C_1 = 0 \text{ or } p = 4/(n-2), \quad (4.9a)$$

$$Y_2 = \xi^{2-n}(\xi/(n-2) \partial/\partial \xi - A \partial/\partial A), \quad \text{for } p = 2(3-n)/(n-2), \quad (4.9b)$$

$$Y_3 = 2((C_1^2/k)\xi^{1/3} - \xi) \partial/\partial \xi - A \partial/\partial A, \quad \text{for } C_1 \neq 0, p = -4, n = 4/3, \quad (4.9c)$$

$$Y_4 = 2(k - C_1^2 \xi^2) \xi \partial / \partial \xi + (k - 4C_1^2 \xi^2) A \partial / \partial A, \quad \text{for } C_1 \neq 0, p = -4, n = 0, \quad (4.9d)$$

$$Y_5 = \xi^3 \partial / \partial \xi - 6(\xi^2 A - 16/k) \partial / \partial A, \quad \text{for } C_1 = 0, p = 1, n = 16, \quad (4.9e)$$

$$Y_6 = 9\xi^{2/3} \partial / \partial \xi - (12\xi^{-1/3} A + (4/k)\xi^{-7/3}) \partial / \partial A, \quad \text{for } C_1 = 0, p = 1, n = 13/3. \quad (4.9f)$$

Since the point symmetries (4.9a) and (4.9b) are the only ones also admitted by the polar system (4.1), the remaining 4 point symmetries (4.9c)–(4.9f) thus describe hidden conditional symmetries arising only for level-set solutions of the polar system (4.1) in the case $\nu = 0$, $C_1 \neq 0$ or $p \neq -1$.

We now carry out the two reduction methods (cf. Secs. III A and III B) to obtain explicit solutions $U(\xi)$ for the translation-group ODE (2.51). The first reduction method can be applied only in the case $\nu = 0$, using the scaling and dilation symmetries (2.59) and (2.60). The second reduction method is applicable to the case $\nu = 0$ by using the hidden conditional symmetries (4.9c)–(4.9f), and the case $\nu \neq 0$ by using the linearization which holds when $C_1 = 0$, $p = -1$.

A. Scaling-symmetry quadratures for the translation-group ODE

The canonical coordinates of the scaling symmetry (2.59) are

$$z = \ln \xi, \quad V = \xi^{2/p} U \quad (4.10)$$

with $\nu = 0$. Hence the translation-group ODE (2.51) gets transformed into

$$V'' + (n - 2 - 4/p)V' + (2(2 - n)/p + 4/p^2)V + k|V|^p V = 0, \quad (4.11)$$

which is equivalent to the ODE (2.61) describing static solutions of the radial Schrödinger equation. To apply the reduction Lemma 2 to ODE (4.11), we note $a = 1$, $b = n - 2 - 4/p$, $c = 2(2 - n)/p + (2/p)^2$.

Then from Lemma 2(i) we obtain the invariant solution

$$V = (2(n - 2 - 2/p)/(kp))^{1/p} \exp(i\Phi), \quad \Phi = \text{const.} \quad (4.12)$$

This yields

$$U = (2(n - 2 - 2/p)/(kp))^{1/p} \xi^{-2/p} \exp(i\Phi), \quad \Phi = \text{const.}, \quad (4.13)$$

which is the scaling-invariant solution of ODE (2.61).

Next, by Lemma 2(ii), we can obtain the general solution for $V(z)$ in the case when $n - 2 - 4/p = 0$ under which the ODE (4.11) becomes

$$V'' - (2/p)^2 V + k|V|^p V = 0, \quad p = 4/(n - 2). \quad (4.14)$$

The quadratures (3.16) in polar variables $V = A \exp(i\Phi)$ are then given by

$$\int \frac{dA}{\sqrt{H(A)}} = \pm z + C_3, \quad \Phi = C_2 \int \frac{dA}{A^2 \sqrt{H(A)}} + C_4, \quad (4.15)$$

(after renaming constants) where

$$H(A) = C_1 - C_2^2 A^{-2} + (2/p)^2 A^2 - 2k \int A^{1+p} dA. \quad (4.16)$$

These integrals (4.15) cannot be evaluated generally to obtain explicit solutions for $A(z)$ and $\Phi(z)$ in terms of elementary functions. Some special cases where explicit solutions can be derived are possible if we make a change of variables

$$\int \frac{dA}{\sqrt{H(A)}} = (1/s) \int \frac{dy}{\sqrt{H_1(y)}}, \quad y = A^s \quad (4.17)$$

and, when $C_2 \neq 0$,

$$\int \frac{dA}{A^2 \sqrt{H(A)}} = (1/\bar{s}) \int \frac{dy}{\sqrt{H_2(y)}}, \quad y = A^{\bar{s}} \quad (4.18)$$

so that the respective expressions

$$H_1(y) = y^{2-2/s} H(y^{1/s}) = C_1 y^{2-2/s} - C_2^2 y^{2-4/s} + (2/p)^2 y^2 - \tilde{k} y^{2+p/s} \quad (4.19)$$

and

$$H_2(y) = y^{2+2/\tilde{s}} H(y^{1/\tilde{s}}) = C_1 y^{2+2/\tilde{s}} - C_2^2 y^2 + (2/p)^2 y^{2+4/\tilde{s}} - \tilde{k} y^{2+(4+p)/\tilde{s}} \quad \text{for } C_2 \neq 0 \quad (4.20)$$

are quadratic polynomials in y for some values of s and \tilde{s} , depending on p , C_1 , C_2 , where

$$\tilde{k} = 2k/(p+2), \quad p = 4/(n-2) \neq -2. \quad (4.21)$$

The required conditions from expression (4.19) consist of

$$2 + p/s = 0, 1, 2; \quad (4.22a)$$

$$2 - 2/s = 0, 1, 2 \quad \text{if } C_1 \neq 0; \quad (4.22b)$$

$$2 - 4/s = 0, 1, 2 \quad \text{if } C_2 \neq 0. \quad (4.22c)$$

From expression (4.20), the required conditions are given by

$$2 + 4/\tilde{s} = 0, 1, 2 \text{ and } 2 + (4+p)/\tilde{s} = 0, 1, 2 \quad \text{if } C_2 \neq 0; \quad (4.22d)$$

$$2 + 2/\tilde{s} = 0, 1, 2 \quad \text{if } C_1 \neq 0, C_2 \neq 0. \quad (4.22e)$$

It is straightforward to solve these conditions for p , s , \tilde{s} . For each value found for p , only a single value each for s and \tilde{s} is needed to evaluate the integrals (4.17) and (4.18). We thus find

$$C_1 = 0 \text{ and } C_2 = 0 : \quad p \neq -2, s = -p/2; \quad (4.23)$$

$$C_1 \neq 0 \text{ and } C_2 = 0 : \quad p = -1, s = 1; \quad p = -4, s = 2; \quad (4.24)$$

$$C_1 = 0 \text{ and } C_2 \neq 0 : \quad p = -4, s = 2, \tilde{s} = -2; \quad (4.25)$$

$$C_1 \neq 0 \text{ and } C_2 \neq 0 : \quad p = -4, s = 2, \tilde{s} = -2; \quad (4.26)$$

with $n = 2 + 4/p$. By scaling and shifting y , we can then match the integrals (4.17) and (4.18) to one of the forms

$$\int \frac{dy}{\sqrt{y^2-1}} = \operatorname{arccosh}(y), \quad \int \frac{dy}{\sqrt{y^2+1}} = \operatorname{arcsinh}(y), \quad (4.27)$$

$$\int \frac{dy}{\sqrt{y^2}} = \ln|y|, \quad (4.28)$$

$$\int \frac{dy}{\sqrt{1-y^2}} = \operatorname{arcsin}(y), \quad (4.29)$$

$$\int \frac{dy}{\sqrt{y}} = 2\sqrt{y}, \quad (4.30)$$

$$\int dy = y. \quad (4.31)$$

More generally, a similar method can be used to obtain additional explicit solutions for $A(z)$ and $\Phi(z)$ by requiring that the expressions (4.19) and (4.20) are either quartic polynomials in y , which will yield elliptic functions, or squares of quartic polynomials in y , which will yield elementary functions. (These solutions will be worked out in Ref. 12.)

1. Quadrature for $p \neq -2$

From case (4.23), we have $s = -p/2$, $C_1 = 0$, and $C_2 = 0$. The quadratures (3.16) are thus given by

$$\Phi = C_4 \quad (4.32)$$

and

$$\pm z + C_3 = (-2/p) \int \frac{dy}{\sqrt{(4/p^2)y^2 - 2k/(p+2)}}, \quad y = A^{-p/2}, \quad (4.33)$$

which can be matched to the form (4.27). This yields

$$A^{-p/2} = p(2(p+2)/k)^{-1/2} \cosh(\pm z + C_3), \quad \Phi = C_4, \quad k/(p+2) > 0, \quad (4.34a)$$

$$A^{-p/2} = p(-2(p+2)/k)^{-1/2} \sinh(\pm z + C_3), \quad \Phi = C_4, \quad k/(p+2) < 0. \quad (4.34b)$$

Hence we obtain two solutions of ODE (4.14) for $p \neq -2$:

$$V = (2(p+2)/(kp^2))^{1/p} (\cosh(\pm z + C_3))^{-2/p} \exp(iC_4), \quad k/(p+2) > 0, \quad (4.35a)$$

$$V = (-2(p+2)/(kp^2))^{1/p} (\sinh(\pm z + C_3))^{-2/p} \exp(iC_4), \quad k/(p+2) < 0. \quad (4.35b)$$

These solutions (4.35a) and (4.35b) can be merged after we change variables (4.10), giving the following result.

Proposition 3. For $p \neq -2$, the translation-group ODE (2.61) has a solution

$$U = (\pm 8(p+2)/(kp^2))^{1/p} (\tilde{C}_3 \xi^2 \pm 1/\tilde{C}_3)^{-2/p} \exp(iC_4), \quad (4.36)$$

$$n = 2 + 4/p, \quad \pm k/(p+2) > 0,$$

with real constants \tilde{C}_3, C_4 .

2. Quadrature for $p = -1$

From case (4.24), we have $s = 1$, $C_1 \neq 0$ and $C_2 = 0$. The quadratures (3.16) are thus given by

$$\Phi = C_4 \quad (4.37)$$

and

$$\pm z + C_3 = \int \frac{dy}{\sqrt{4(y - k/4)^2 + C_1 - k^2/4}}, \quad y = A, \quad (4.38)$$

which can be matched to the forms (4.27) and (4.28). This yields

$$A = \frac{1}{4} (k + (4C_1 - k^2)^{1/2} \sinh(2(\pm z + C_3))), \quad \Phi = C_4, \quad C_1 > k^2/4, \quad (4.39a)$$

$$A = \frac{1}{4} (k + (k^2 - 4C_1)^{1/2} \cosh(2(\pm z + C_3))), \quad \Phi = C_4, \quad C_1 < k^2/4, \quad (4.39b)$$

$$A = \frac{1}{4} (k + 4 \exp(2(\pm z + C_3))), \quad \Phi = C_4, \quad C_1 = k^2/4, \quad (4.39c)$$

$$A = \frac{1}{4} (k - 4 \exp(2(\pm z + C_3))), \quad \Phi = C_4, \quad C_1 = k^2/4. \quad (4.39d)$$

Hence we obtain four solutions of ODE (4.14) for $p = -1$:

$$V = \frac{1}{4} (k + (4C_1 - k^2)^{1/2} \sinh(2(\pm z + C_3))) \exp(iC_4), \quad C_1 > k^2/4, \quad (4.40a)$$

$$V = \frac{1}{4} (k + (k^2 - 4C_1)^{1/2} \cosh(2(\pm z + C_3))) \exp(iC_4), \quad C_1 < k^2/4, \quad (4.40b)$$

$$V = \frac{1}{4} (k \pm 4 \exp(2(z + C_3))) \exp(iC_4), \quad C_1 = k^2/4, \quad (4.40c)$$

$$V = \frac{1}{4} (k \pm 4 \exp(2(-z + C_3))) \exp(iC_4), \quad C_1 = k^2/4. \quad (4.40d)$$

These solutions (4.40a)–(4.40d) give the following result after we change variables (4.10).

Proposition 4. For $p = -1$, the translation-group ODE (2.61) has solutions

$$U = \frac{1}{4} (k\xi^2 + (\pm(C_1 - k^2/4))^{1/2} (\tilde{C}_3 \xi^4 \mp 1/\tilde{C}_3)) \exp(iC_4), \quad (4.41)$$

$$C_1 \neq k^2/4, \quad n = -2,$$

$$U = \frac{1}{4} (k\xi^2 + \tilde{C}_3 \xi^4) \exp(iC_4), \quad n = -2, \quad (4.42)$$

$$U = \frac{1}{4} (k\xi^2 + \tilde{C}_3) \exp(iC_4), \quad n = -2, \quad (4.43)$$

with real constants C_1, \tilde{C}_3, C_4 .

3. Quadrature for $p = -4$

From cases (4.24)–(4.26), we have $s = 2$, $C_1 \neq 0$ and $C_2 = 0$; $s = 2$, $\tilde{s} = -2$, and $C_2 \neq 0$.

In the first subcase, the quadratures (3.16) are given by

$$\Phi = C_4 \quad (4.44)$$

and

$$\pm z + C_3 = \int \frac{dy}{\sqrt{(y + 2C_1)^2 + 4(k - C_1^2)}}, \quad y = A^2. \quad (4.45)$$

Quadrature (4.45) can be matched to the forms (4.27) and (4.28). This yields

$$A^2 = -2C_1 + 2(k - C_1^2)^{1/2} \sinh(\pm z + C_3), \quad \Phi = C_4, \quad C_1^2 < k, \quad (4.46a)$$

$$A^2 = -2C_1 + 2(C_1^2 - k)^{1/2} \cosh(\pm z + C_3), \quad \Phi = C_4, \quad C_1^2 > k, \quad (4.46b)$$

$$A^2 = -2C_1 + \exp(\pm z + C_3), \quad \Phi = C_4, \quad C_1^2 = k, \quad (4.46c)$$

$$A^2 = -2C_1 - \exp(\pm z + C_3), \quad \Phi = C_4, \quad C_1^2 = k, C_1 < 0. \quad (4.46d)$$

Hence we obtain four solutions of ODE (4.14) for $p = -4$:

$$V = (-2C_1 + 2(k - C_1^2)^{1/2} \sinh(\pm z + C_3))^{1/2} \exp(iC_4), \quad C_1^2 < k, \quad (4.47a)$$

$$V = (-2C_1 + 2(C_1^2 - k)^{1/2} \cosh(\pm z + C_3))^{1/2} \exp(iC_4), \quad C_1^2 > k, \quad (4.47b)$$

$$V = (-2C_1 + \exp(\pm z + C_3))^{1/2} \exp(iC_4), \quad C_1^2 = k, \quad (4.47c)$$

$$V = (-2C_1 - \exp(\pm z + C_3))^{1/2} \exp(iC_4), \quad C_1^2 = k, C_1 < 0. \quad (4.47d)$$

In the second subcase, the quadratures (3.16) are given by

$$\pm z + C_3 = \int \frac{dy}{\sqrt{(y + 2C_1)^2 + 4(k - C_1^2 - C_2^2)}}, \quad y = A^2 \quad (4.48)$$

and

$$\Phi = C_4 - C_2 \int \frac{dy}{\sqrt{4(k - C_2^2)y^2 + 4C_1y + 1}}, \quad y = A^{-2}. \quad (4.49)$$

Quadrature (4.48) can be matched to the forms (4.27) and (4.28), which yields

$$A^2 = -2C_1 + 2(k - C_1^2 - C_2^2)^{1/2} \sinh(\pm z + C_3), \quad C_1^2 + C_2^2 < k, \quad (4.50a)$$

$$A^2 = -2C_1 + 2(C_2^2 + C_1^2 - k)^{1/2} \cosh(\pm z + C_3), \quad C_1^2 + C_2^2 > k, \quad (4.50b)$$

$$A^2 = -2C_1 + \exp(\pm z + C_3), \quad C_1^2 + C_2^2 = k, \quad (4.50c)$$

$$A^2 = -2C_1 - \exp(\pm z + C_3), \quad C_1^2 + C_2^2 = k, C_1 < 0. \quad (4.50d)$$

Quadrature (4.49) can be matched to all of the forms (4.27)–(4.31). This yields

$$\begin{aligned} \Phi &= C_4 + \frac{1}{2}(k/C_2^2 - 1)^{-1/2} \operatorname{arcsinh} \left((k - C_2^2 - C_1^2)^{-1/2} (C_1 + 2(k - C_2^2)A^{-2}) \right), \\ &C_1^2 + C_2^2 < k, \end{aligned} \quad (4.51a)$$

$$\begin{aligned} \Phi &= C_4 - \frac{1}{2}(k/C_2^2 - 1)^{-1/2} \operatorname{arccosh} \left((C_2^2 + C_1^2 - k)^{-1/2} (C_1 + 2(k - C_2^2)A^{-2}) \right), \\ &C_2^2 < k < C_1^2 + C_2^2, \end{aligned} \quad (4.51b)$$

$$\begin{aligned} \Phi &= C_4 - \frac{1}{2}(1 - k/C_2^2)^{-1/2} \arcsin \left((C_2^2 + C_1^2 - k)^{-1/2} (C_1 - 2(C_2^2 - k)A^{-2}) \right), \\ &C_2^2 > k, \end{aligned} \quad (4.51c)$$

$$\Phi = C_4 - (C_2/(2C_1)) \ln |C_1 A^{-2} + 1/2|, \quad C_1^2 + C_2^2 = k, \quad (4.51d)$$

$$\Phi = C_4 - (C_2/C_1)(C_1 A^{-2} + 1/4)^{1/2}, \quad C_2^2 = k, \quad (4.51e)$$

$$\Phi = C_4 - C_2 A^{-2}, \quad C_2^2 = k, C_1 = 0. \quad (4.51f)$$

Hence we obtain 7 more solutions of ODE (4.14) for $p = -4$:

$$\begin{aligned} V &= (-2C_1 + 2(k - C_1^2 - C_2^2)^{1/2} \sinh(\pm z + C_3))^{1/2} \\ &\times \exp \left(iC_4 - (i/2)(k/C_2^2 - 1)^{-1/2} \operatorname{arcsinh} \left(\frac{(k - C_1^2 - C_2^2)^{1/2} + C_1 \sinh(\pm z + C_3)}{-C_1 + (k - C_1^2 - C_2^2)^{1/2} \sinh(\pm z + C_3)} \right) \right), \\ &C_1^2 + C_2^2 < k, \end{aligned} \quad (4.52a)$$

$$\begin{aligned} V &= (-2C_1 + 2(C_2^2 + C_1^2 - k)^{1/2} \cosh(\pm z + C_3))^{1/2} \\ &\times \exp \left(iC_4 - (i/2)(k/C_2^2 - 1)^{-1/2} \operatorname{arccosh} \left(\frac{(C_2^2 + C_1^2 - k)^{1/2} - C_1 \cosh(\pm z + C_3)}{C_1 - (C_2^2 + C_1^2 - k)^{1/2} \cosh(\pm z + C_3)} \right) \right), \\ &C_2^2 < k < C_1^2 + C_2^2, \end{aligned} \quad (4.52b)$$

$$\begin{aligned} V &= (-2C_1 + 2(C_2^2 + C_1^2 - k)^{1/2} \cosh(\pm z + C_3))^{1/2} \\ &\times \exp \left(iC_4 - (i/2)(1 - k/C_2^2)^{-1/2} \arcsin \left(\frac{(C_2^2 + C_1^2 - k)^{1/2} - C_1 \cosh(\pm z + C_3)}{-C_1 + (C_2^2 + C_1^2 - k)^{1/2} \cosh(\pm z + C_3)} \right) \right), \\ &C_2^2 > k, \end{aligned} \quad (4.52c)$$

$$\begin{aligned} V &= (-2C_1 + \exp(\pm z + C_3))^{1/2} \exp \left(i\tilde{C}_4 + i(C_2/(2C_1)) \ln |2C_1 \exp(\mp z - C_3) - 1| \right), \\ &C_1^2 + C_2^2 = k, \end{aligned} \quad (4.52d)$$

$$V = (-2C_1 - \exp(\pm z + C_3))^{1/2} \exp(i\tilde{C}_4 + i(C_2/(2C_1)) \ln |2C_1 \exp(\mp z - C_3) + 1|), \quad (4.52e)$$

$$C_1^2 + C_2^2 = k,$$

$$V = (-2C_1 + 2|C_1| \cosh(\pm z + C_3))^{1/2} \exp\left(iC_4 - i(k^{1/2}/(2C_1)) \left(\frac{C_1 + |C_1| \cosh(\pm z + C_3)}{-C_1 + |C_1| \cosh(\pm z + C_3)}\right)^{1/2}\right),$$

$$C_2^2 = k, \quad (4.52f)$$

$$V = \exp(\frac{1}{2}(\pm z + C_3)) \exp(iC_4 - ik^{1/2} \exp(\mp z - C_3)), \quad C_2^2 = k, C_1 = 0. \quad (4.52g)$$

After we change variables (4.10), these solutions (4.47a)–(4.47d) and (4.52a)–(4.52g) give the following result.

Proposition 5. For $p = -4$, the translation-group ODE (2.61) has solutions

$$U = (-2C_1\xi + (\pm(k - C_1^2))^{1/2}(\tilde{C}_3\xi^2 \mp 1/\tilde{C}_3))^{1/2} \exp(iC_4), \quad C_1^2 \neq k, \quad n = 1, \quad (4.53)$$

$$U = (\pm 2k^{1/2}\xi + \tilde{C}_3\xi^2)^{1/2} \exp(iC_4), \quad k > 0, \quad n = 1, \quad (4.54)$$

$$U = (\pm 2k^{1/2}\xi + \tilde{C}_3)^{1/2} \exp(iC_4), \quad k > 0, \quad n = 1, \quad (4.55)$$

$$U = (-2C_1\xi + (k - C_1^2 - C_2^2)^{1/2}(\tilde{C}_3\xi^2 - 1/\tilde{C}_3))^{1/2} \\ \times \exp\left(iC_4 - \frac{1}{2}i(k/C_2^2 - 1)^{-1/2} \operatorname{arcsinh}\left(\frac{2(k - C_1^2 - C_2^2)^{1/2}\xi + C_1(\tilde{C}_3\xi^2 - 1/\tilde{C}_3)}{-2C_1\xi + (k - C_1^2 - C_2^2)^{1/2}(\tilde{C}_3\xi^2 - 1/\tilde{C}_3)}\right)\right),$$

$$C_1^2 + C_2^2 < k, \quad n = 1, \quad (4.56)$$

$$U = (-2C_1\xi + (C_2^2 + C_1^2 - k)^{1/2}(\tilde{C}_3\xi^2 + 1/\tilde{C}_3))^{1/2} \\ \times \exp\left(iC_4 - \frac{1}{2}i(k/C_2^2 - 1)^{-1/2} \operatorname{arccosh}\left(\frac{2(C_2^2 + C_1^2 - k)^{1/2}\xi - C_1(\tilde{C}_3\xi^2 + 1/\tilde{C}_3)}{2C_1\xi - (C_2^2 + C_1^2 - k)^{1/2}(\tilde{C}_3\xi^2 + 1/\tilde{C}_3)}\right)\right),$$

$$C_2^2 < k < C_1^2 + C_2^2, \quad n = 1, \quad (4.57)$$

$$U = (-2C_1\xi + (C_2^2 + C_1^2 - k)^{1/2}(\tilde{C}_3\xi^2 + 1/\tilde{C}_3))^{1/2} \\ \times \exp\left(iC_4 - \frac{1}{2}i(1 - k/C_2^2)^{-1/2} \arcsin\left(\frac{2(C_2^2 + C_1^2 - k)^{1/2}\xi - C_1(\tilde{C}_3\xi^2 + 1/\tilde{C}_3)}{-2C_1\xi + (C_2^2 + C_1^2 - k)^{1/2}(\tilde{C}_3\xi^2 + 1/\tilde{C}_3)}\right)\right),$$

$$C_2^2 > k, \quad n = 1, \quad (4.58)$$

$$U = (-2C_1\xi + \tilde{C}_3\xi^2)^{1/2} \exp(i\tilde{C}_4 + i(C_2/(2C_1)) \ln |2C_1\xi^{-1} - \tilde{C}_3|), \quad C_1^2 + C_2^2 = k, \quad n = 1, \quad (4.59)$$

$$U = (-2C_1\xi + \tilde{C}_3)^{1/2} \exp(i\tilde{C}_4 + i(C_2/(2C_1)) \ln |2C_1\xi - \tilde{C}_3|), \quad C_1^2 + C_2^2 = k, \quad n = 1, \quad (4.60)$$

$$U = (C_1/\tilde{C}_3)^{1/2}(\tilde{C}_3\xi - 1) \exp\left(iC_4 - \frac{1}{2}i(k^{1/2}/C_1) \left(\frac{\tilde{C}_3\xi + 1}{\tilde{C}_3\xi - 1}\right)\right), \quad k > 0, \quad n = 1, \quad (4.61)$$

$$U = \tilde{C}_3^{1/2} \xi \exp(iC_4 - i(k^{1/2}/\tilde{C}_3)\xi^{-1}), \quad k > 0, \quad n = 1, \quad (4.62)$$

$$U = \tilde{C}_3^{1/2} \exp(iC_4 - i(k^{1/2}/\tilde{C}_3)\xi), \quad k > 0, \quad n = 1, \quad (4.63)$$

with real constants $C_1, C_2, \tilde{C}_3, C_4, \tilde{C}_4$.

B. Dilation-symmetry quadratures for the translation-group ODE

The canonical coordinates of the dilation symmetry (2.60) are

$$z = \xi^{n-2}, \quad V = \xi^{n-2}U \quad (4.64)$$

with $v = 0$. In this case the translation-group ODE (2.61) gets transformed into

$$V'' + k(1 + p/2)^2 |V|^p V = 0, \quad p = 2(3 - n)/(n - 2), \quad n \neq 2, 3. \quad (4.65)$$

To apply Lemma 2 to this ODE (4.65), we note $a = 1/(1 + p/2)^2, b = 0, c = 0$.

From Lemma 2(i) we see that the invariant solution is trivial, since the only root of F is $|V| = 0$. Thus ODE (2.61) has no non-trivial dilation-invariant solution.

Next, we can obtain the general solution for $V(z)$ from the quadratures (3.16), since the conditions of Lemma 2(ii) are directly satisfied for the ODE (4.65). In polar variables $V = A \exp(i\Phi)$, these quadratures are given by

$$\int \frac{dA}{\sqrt{H(A)}} = \pm z + C_3, \quad \Phi = C_2 \int \frac{dA}{A^2 \sqrt{H(A)}} + C_4, \quad (4.66)$$

where

$$H(A) = C_1 - C_2^2 A^{-2} - 2k(1 + p/2)^2 \int A^{1+p} dA. \quad (4.67)$$

These integrals (4.66) cannot be evaluated generally to obtain explicit solutions for $A(z)$ and $\Phi(z)$ in terms of elementary functions. Some special cases where explicit solutions can be derived are possible if we change variables (4.17) and (4.18) to make the expressions

$$H_1(y) = y^{2-2/s} H(y^{1/s}) = C_1 y^{2-2/s} - C_2^2 y^{2-4/s} - \tilde{k} y^{2+p/s}, \quad (4.68)$$

$$H_2(y) = y^{2+2/\tilde{s}} H(y^{1/\tilde{s}}) = C_1 y^{2+2/\tilde{s}} - C_2^2 y^2 - \tilde{k} y^{2+(4+p)/\tilde{s}} \quad \text{for } C_2 \neq 0 \quad (4.69)$$

into quadratic polynomials in y for some values of s and \tilde{s} , depending on p, C_1, C_2 , where

$$\tilde{k} = 2k/(p + 2), \quad p = 2(3 - n)/(n - 2) \neq -2. \quad (4.70)$$

The required conditions from expressions (4.68) and (4.69) are, respectively, given by (4.22a)–(4.22c), (4.22e), and

$$2 + (4 + p)/\tilde{s} = 0, 1, 2 \quad \text{if } C_2 \neq 0. \quad (4.71)$$

Solving these conditions for p, s, \tilde{s} in the same way as for the previous case, we find

$$C_1 = 0 \text{ and } C_2 = 0: \quad p \neq -2, s = -p/2; \quad (4.72)$$

$$C_1 \neq 0 \text{ and } C_2 = 0: \quad p = -1, s = 1; \quad p = -4, s = 2; \quad (4.73)$$

$$C_1 = 0 \text{ and } C_2 \neq 0: \quad p = -4, s = 2, \tilde{s} = -1; \quad p = -8, s = 4, \tilde{s} = 2; \quad (4.74)$$

$$C_1 \neq 0 \text{ and } C_2 \neq 0: \quad p = -4, s = 2, \tilde{s} = -1; \quad (4.75)$$

with $n = 2(p + 3)/(p + 2)$.

We can then match the integrals (4.17) and (4.18) to one of the forms (4.27)–(4.31) after scaling and shifting y .

As noted previously, additional explicit solutions for $A(z)$ and $\Phi(z)$ can be derived by requiring that the expressions (4.68) and (4.69) are either quartic polynomials in y , which will yield elliptic functions, or squares of quartic polynomials in y , which will yield elementary functions. (These solutions will be worked out in Ref. 12.)

1. Quadrature for $p \neq -2$

From case (4.72), we have $s = -p/2$, $C_1 = 0$, and $C_2 = 0$. The quadratures (3.16) are thus given by

$$\Phi = C_4 \quad (4.76)$$

and

$$\pm z + C_3 = -(2/p) \int \frac{dy}{\sqrt{-k(p+2)/2}}, \quad y = A^{-p/2}, \quad (4.77)$$

which can be matched to the form (4.31). This yields

$$A^{-p/2} = (-kp^2(p+2)/8)^{1/2}(\pm z + C_3), \quad \Phi = C_4, \quad k(p+2) < 0. \quad (4.78)$$

Hence we obtain a solution of ODE (4.65) for $p \neq -2$:

$$V = (-kp^2(p+2)/8)^{-1/p}(\pm z + C_3)^{-2/p} \exp(iC_4), \quad k(p+2) < 0. \quad (4.79)$$

This solution (4.79) gives the following result after we change variables (4.64).

Proposition 6. For $p \neq -2$, the translation-group ODE (2.61) has a solution

$$U = (-kp^2(p+2)/8)^{-1/p}(\xi + \tilde{C}_3\xi^{3-n})^{-2/p} \exp(iC_4), \quad (4.80)$$

$$k(p+2) < 0, \quad n = 2(p+3)/(p+2),$$

with real constants \tilde{C}_3, C_4 .

2. Quadrature for $p = -1$

From case (4.73), we have $s = 1$, $C_1 \neq 0$, and $C_2 = 0$. The quadratures (3.16) are thus given by

$$\Phi = C_4 \quad (4.81)$$

and

$$\pm z + C_3 = \int \frac{dy}{\sqrt{C_1 - (k/2)y}}, \quad y = A, \quad (4.82)$$

which can be matched to the form (4.30). This yields

$$A = (2C_1/k) - (k/8)(\pm z + C_3)^2, \quad \Phi = C_4. \quad (4.83)$$

Hence we obtain a solution of ODE (4.65) for $p = -1$:

$$V = ((2C_1/k) - (k/8)(\pm z + C_3)^2) \exp(iC_4). \quad (4.84)$$

This solution (4.84) gives the following result after we change variables (4.64).

Proposition 7. For $p = -1$, the translation-group ODE (2.61) has a solution

$$U = ((2C_1/k)\xi^{-2} - (k/8)(\xi + \tilde{C}_3/\xi)^2) \exp(iC_4), \quad n = 4, \quad (4.85)$$

with real constants C_1, \tilde{C}_3, C_4 .

3. Quadrature for $p = -4$

From cases (4.73)–(4.75), we have $s = 2$, $\tilde{s} = -1$, and $C_2 \neq 0$.

The quadratures (3.16) are thus given by

$$\pm z + C_3 = (1/2) \int \frac{dy}{\sqrt{C_1 y + k - C_2^2}}, \quad y = A^2, \quad (4.86)$$

which can be matched to the forms (4.30) and (4.31), and

$$\Phi = C_4 - C_2 \int \frac{dy}{\sqrt{C_1 + (k - C_2^2)y^2}}, \quad y = A^{-1}, \quad (4.87)$$

which can be matched to all of the forms (4.27)–(4.29). This yields

$$A^2 = (C_2^2 - k)/C_1 + C_1(\pm z + C_3)^2, \quad C_1 \neq 0, \quad (4.88a)$$

$$A^2 = 2(k - C_2^2)^{1/2}(\pm z + C_3), \quad C_2^2 < k, C_1 = 0, \quad (4.88b)$$

and

$$\Phi = C_4 - C_2(k - C_2^2)^{-1/2} \operatorname{arcsinh}(((k - C_2^2)/C_1)^{1/2} A^{-1}), \quad C_2^2 < k, C_1 > 0, \quad (4.89a)$$

$$\Phi = C_4 - C_2(C_2^2 - k)^{-1/2} \arcsin(((C_2^2 - k)/C_1)^{1/2} A^{-1}), \quad C_2^2 > k, C_1 > 0, \quad (4.89b)$$

$$\Phi = C_4 - C_2(k - C_2^2)^{-1/2} \operatorname{arccosh}(((C_2^2 - k)/C_1)^{1/2} A^{-1}), \quad C_2^2 < k, C_1 < 0, \quad (4.89c)$$

$$\Phi = C_4 - (C_2/C_1^{1/2})A^{-1}, \quad C_2^2 = k, C_1 > 0, \quad (4.89d)$$

$$\Phi = C_4 - C_2(k - C_2^2)^{-1/2} \ln |A^{-1}|, \quad C_2^2 < k, C_1 = 0. \quad (4.89e)$$

Hence we obtain 5 more solutions of ODE (4.14) for $p = -4$:

$$V = ((C_2^2 - k)/C_1 + C_1(\pm z + C_3)^2)^{1/2} \times \exp \left(iC_4 - iC_2(k - C_2^2)^{-1/2} \operatorname{arcsinh}((-1 + C_1^2(\pm z + C_3)^2/(k - C_2^2))^{-1/2}) \right), \quad (4.90a)$$

$$C_2^2 < k, C_1 > 0,$$

$$V = ((C_2^2 - k)/C_1 + C_1(\pm z + C_3)^2)^{1/2} \times \exp \left(iC_4 - iC_2(C_2^2 - k)^{-1/2} \arcsin \left((1 + C_1^2(\pm z + C_3)^2/(C_2^2 - k))^{-1/2} \right) \right), \quad (4.90b)$$

$$C_2^2 > k, C_1 > 0,$$

$$V = ((C_2^2 - k)/C_1 + C_1(\pm z + C_3)^2)^{1/2} \times \exp \left(iC_4 - iC_2(k - C_2^2)^{-1/2} \operatorname{arccosh} \left((1 + C_1^2(\pm z + C_3)^2/(C_2^2 - k))^{-1/2} \right) \right), \quad (4.90c)$$

$$C_2^2 < k, C_1 < 0,$$

$$V = C_1^{1/2} |\pm z + C_3| \exp \left(iC_4 \mp i(k^{1/2}/C_1) |\pm z + C_3|^{-1} \right), \quad C_2^2 = k, C_1 > 0, \quad (4.90d)$$

$$V = (4(k - C_2^2))^{1/4} (\pm z + C_3)^{1/2} \exp \left(i\tilde{C}_4 + \frac{1}{2} iC_2(k - C_2^2)^{-1/2} \ln |\pm z + C_3| \right), \quad (4.90e)$$

$$C_2^2 < k, C_1 = 0.$$

After we change variables (4.64), these solutions (4.90a)–(4.90e) give the following result.

Proposition 8. For $p = -4$, the translation-group ODE (2.61) has solutions

$$U = ((C_2^2 - k)\xi^2/C_1 + C_1(1 + \tilde{C}_3\xi)^2)^{1/2} \times \exp\left(iC_4 - iC_2(k - C_2^2)^{-1/2} \operatorname{arcsinh}\left((-1 + C_1^2(\xi^{-1} + \tilde{C}_3)^2/(k - C_2^2))^{-1/2}\right)\right), \quad (4.91)$$

$$C_2^2 < k, \quad C_1 > 0, \quad n = 1,$$

$$U = ((C_2^2 - k)\xi^2/C_1 + C_1(1 + \tilde{C}_3\xi)^2)^{1/2} \times \exp\left(iC_4 - iC_2(C_2^2 - k)^{-1/2} \arcsin\left((1 + C_1^2(\xi^{-1} + \tilde{C}_3)^2/(C_2^2 - k))^{-1/2}\right)\right), \quad (4.92)$$

$$C_2^2 > k, \quad C_1 > 0, \quad n = 1,$$

$$U = ((C_2^2 - k)\xi^2/C_1 + C_1(1 + \tilde{C}_3\xi)^2)^{1/2} \times \exp\left(iC_4 - iC_2(k - C_2^2)^{-1/2} \operatorname{arccosh}\left((1 + C_1^2(\xi^{-1} + \tilde{C}_3)^2/(C_2^2 - k))^{-1/2}\right)\right), \quad (4.93)$$

$$C_2^2 < k, \quad C_1 < 0, \quad n = 1,$$

$$U = (4(k - C_2^2))^{1/4} (C_3\xi^2 \pm \xi)^{1/2} \exp\left(i\tilde{C}_4 + \frac{1}{2}iC_2(k - C_2^2)^{-1/2} \ln|C_3 \pm \xi^{-1}|\right), \quad (4.94)$$

$$C_2^2 < k, \quad n = 1,$$

$$U = C_1^{1/2}|1 + \tilde{C}_3\xi| \exp\left(iC_4 \mp i(k^{1/2}/C_1)|\xi^{-1} + \tilde{C}_3|^{-1}\right), \quad C_1 > 0, \quad n = 1, \quad (4.95)$$

with real constants $C_1, C_2, \tilde{C}_3, C_4, \tilde{C}_4$.

4. Quadrature for $p = -8$

From case (4.74), we have $s = 4$, $\tilde{s} = 2$, $C_1 = 0$, and $C_2 \neq 0$. The quadratures (3.16) are thus given by

$$\pm z + C_3 = (1/4) \int \frac{dy}{\sqrt{3k - C_2^2 y}}, \quad y = A^4 \quad (4.96)$$

and

$$\Phi = C_4 + (C_2/2) \int \frac{dy}{\sqrt{3k - C_2^2 y^2}}, \quad y = A^2. \quad (4.97)$$

Quadrature (4.96) can be matched to the form (4.30), which yields

$$A^4 = 3k/C_2^2 - 4C_2^2(\pm z + C_3)^2, \quad k > 0. \quad (4.98)$$

Quadrature (4.97) can be matched to the form (4.29). This yields

$$\Phi = C_4 + (1/2) \arcsin\left((3k)^{-1/2} C_2 A^2\right), \quad k > 0. \quad (4.99)$$

Hence we obtain a solution of ODE (4.65) for $p = -8$:

$$V = (3k/C_2^2 - 4C_2^2(\pm z + C_3)^2)^{1/4} \exp\left(iC_4 + \frac{1}{2}i \arcsin((1 - 4C_2^4(\pm z + C_3)^2/(3k))^{1/2})\right), \quad (4.100)$$

$$k > 0.$$

This solution (4.100) gives the following result after we change variables (4.64).

Proposition 9. For $p = -8$, the translation-group ODE (2.61) has a solution

$$U = ((3k/C_2^2)\xi^{4/3} - 4C_2^2(\xi^{1/3} + \tilde{C}_3\xi^{2/3})^2)^{1/4} \times \exp\left(iC_4 + \frac{1}{2}i\arcsin\left((1 - 4C_2^4(\xi^{-1/3} + \tilde{C}_3)^2/(3k))^{1/2}\right)\right), \quad k > 0, \quad n = 5/3, \quad (4.101)$$

with real constants C_2, \tilde{C}_3, C_4 .

C. Conditional-symmetry quadratures for the translation-group ODE

From Proposition 2, the hidden conditional symmetries (4.9c)–(4.9f) with $\nu = 0$ can be naturally split up into two types: symmetries (4.9c) and (4.9d) hold for $p = -4$, $C_1 \neq 0$ and have the form (3.36) to which Lemmas 3 and 4 can be applied; symmetries (4.9e) and (4.9f) hold for $p = 1$, $C_1 = 0$ and have a slightly more general form

$$\tilde{Y} = \zeta(\xi)\partial/\partial\xi + (\Upsilon_1(\xi)A + \Upsilon_0(\xi))\partial/\partial A, \quad (4.102)$$

for which the methods that underlie Lemmas 3 and 4 still can be used.

1. Quadratures for $p = -4$, $C_1 \neq 0$

Point symmetry (4.9c) of the level-set ODE (4.5) has the canonical coordinates

$$z = (-3/4)\ln(C_1^2/k - \xi^{2/3}), \quad B = (C_1^2/k - \xi^{2/3})^{-3/4}A \quad (4.103)$$

with $n = 4/3$, $\nu = 0$. Hence, the Lagrangian (4.6) for ODE (4.5) gets transformed into

$$\hat{L} = \exp(-(2/3)z)((1/2)B'^2 - (1/6)B^2 + 2kB^{-2}), \quad (4.104)$$

yielding the transformed ODE

$$\frac{1}{2}B'' - \frac{1}{3}B' - \frac{1}{6}B - 2kB^{-3} = 0. \quad (4.105)$$

Since the coefficient of B' is non-zero, the reduction shown in Lemma 3 cannot be applied to ODE (4.105). Instead, from Lemma 4, an invariant solution can be derived. This yields

$$B = (-12k)^{1/4}. \quad (4.106)$$

Substituting this solution into expressions (3.51) and (3.52), we obtain

$$A = (-12k)^{1/4}(C_1^2/k - \xi^{2/3})^{3/4}, \quad \Phi = (\sqrt{3}C_1/2)(k\xi^{2/3} - C_1^2)^{-1/2} + C_2, \quad (4.107)$$

which is the corresponding invariant solution of the system (3.26) and (3.27) with $p = -4$, $n = 4/3$, $\nu = 0$.

Point symmetry (4.9d) of the level-set ODE (4.5) has the canonical coordinates

$$z = (1/(4k))\ln(\xi^2/(C_1^2\xi^2 - k)), \quad B = \xi^{-1/2}(C_1^2\xi^2 - k)^{-3/4}A \quad (4.108)$$

with $n = 0$, $\nu = 0$. Hence, the Lagrangian (4.6) for ODE (4.5) gets transformed into

$$\hat{L} = \exp(-2kz)((1/2)B'^2 + (3k^2/2)B^2 - 2B^{-2}). \quad (4.109)$$

This yields the transformed ODE

$$-\frac{1}{2}B'' + kB' + \frac{3}{2}k^2B + 2B^{-3} = 0, \quad (4.110)$$

whose coefficient of B' is non-zero. Again, while the reduction shown in Lemma 3 cannot be applied, an invariant solution to ODE (4.110) can be derived from Lemma 4, yielding

$$B = (-3k^2/4)^{-1/4}. \quad (4.111)$$

Substituting this solution into expressions (3.51) and (3.52), we obtain

$$A = (3k^2/4)^{-1/4} \xi^{1/2} (k - C_1^2 \xi^2)^{3/4}, \quad \Phi = -(\sqrt{3}C_1/2)(k\xi^{-2} - C_1^2)^{-1/2} + C_2, \quad (4.112)$$

which is the corresponding invariant solution of the system (3.26) and (3.27) with $p = -4$, $n = 0$, $\nu = 0$.

These two solutions (4.107) and (4.112) give the following result.

Proposition 10. For $p = -4$, the translation-group ODE (2.61) has solutions

$$U = (-12k)^{1/4} (\tilde{C}_1/k - \xi^{2/3})^{3/4} \exp(i(\sqrt{3}/2)((k/\tilde{C}_1)\xi^{2/3} - 1)^{-1/2} + iC_2), \quad (4.113)$$

$$k < 0, \quad n = 4/3,$$

$$U = (4k/3)^{1/4} \xi^2 (\xi^{-2} - \tilde{C}_1/k)^{3/4} \exp(-i(\sqrt{3}/2)((k/\tilde{C}_1)\xi^{-2} - 1)^{-1/2} + iC_2), \quad (4.114)$$

$$n = 0,$$

with real constants \tilde{C}_1, C_2 .

2. Quadratures for $p = 1, C_1 = 0$

Point symmetry (4.9e) of the level-set ODE (4.5) has the canonical coordinates

$$z = -(1/2)\xi^{-2}, \quad B = \xi^6 A - (24/k)\xi^4 \quad (4.115)$$

with $n = 16$, $C_1 = 0$, $\nu = 0$. Hence the ODE (4.5) gets transformed into

$$B'' + kB^2 = 0. \quad (4.116)$$

This ODE (4.116) admits the integrating factor B' , leading to the quadrature

$$\int \frac{dB}{\sqrt{C_2 - (2k/3)B^3}} = \pm z + C_3, \quad (4.117)$$

which is an elliptic integral when $C_2 \neq 0$. (This solution will be worked out in detail in Ref. 12.) In the case $C_2 = 0$, the integral (4.117) can be evaluated explicitly, yielding

$$B = (-6/k)(\pm z + C_3)^{-2}. \quad (4.118)$$

Hence, from the canonical coordinates (4.115) and the first integral (3.31), we obtain

$$A = (96C_3/k)(C_3\xi^2 \pm 1)/(2C_3\xi^2 \pm 1)^2, \quad \Phi = C_4, \quad (4.119)$$

which is the solution of the system (3.26) and (3.27) with $p = 1$, $n = 16$, $\nu = 0$, and $C_1 = C_2 = 0$. Note that the invariant solution of the ODE (4.116), given by $B' = 0$, is simply $B = 0$. This solution also arises as the special limit $C_3 \rightarrow \infty$ in the solution (4.118). The corresponding invariant solution of system (3.26) and (3.27) is given by

$$A = 24/(k\xi^2), \quad \Phi = C_4, \quad (4.120)$$

which is non-trivial due to the shift term in the change of variables (4.115) for B in terms of A .

Point symmetry (4.9f) of the level-set ODE (4.5) has the canonical coordinates

$$z = (1/3)\xi^{1/3}, \quad B = \xi^{4/3} A - (2/(3k))\xi^{-2/3} \quad (4.121)$$

with $n = 13/3$, $\nu = 0$, $C_1 = 0$. Hence, the ODE (4.5) gets transformed into

$$\frac{1}{81} B'' + kB^2 = 0. \quad (4.122)$$

This ODE (4.122) admits the integrating factor B' , leading to the quadrature

$$\int \frac{dB}{\sqrt{81C_2 - 54kB^3}} = \pm z + C_3. \quad (4.123)$$

When $C_2 \neq 0$, this quadrature is an elliptic integral. (The resulting solution will be worked out in detail in Ref. 12.) In the case $C_2 = 0$, the integral (4.123) can be evaluated explicitly, yielding

$$B = -(2/(27k))(\pm z + C_3)^{-2}. \quad (4.124)$$

The canonical coordinates (4.115) and the first integral (3.31) then give

$$A = (2C_3/k)\xi^{-2}(3C_3 \pm 2\xi^{1/3})/(3C_3 \pm \xi^{1/3})^2, \quad \Phi = C_4, \quad (4.125)$$

which is the solution of the system (3.26) and (3.27) with $p = 1$, $n = 13/3$, $\nu = 0$, and $C_1 = C_2 = 0$. Note that the invariant solution of the ODE (4.122) is again $B = 0$ which also arises as the special limit $C_3 \rightarrow \infty$ in the solution (4.124). The corresponding invariant solution of system (3.26) and (3.27) is given by

$$A = 2/(3k\xi^2), \quad \Phi = C_4. \quad (4.126)$$

The solutions (4.119), (4.120), (4.125), (4.126) give the following result.

Proposition 11. For $p = 1$, the translation-group ODE (2.61) has solutions

$$U = (24/k)\xi^{-2} \exp(iC_4), \quad n = 16, \quad (4.127)$$

$$U = (96/k)(\xi^2 + \tilde{C}_3)(2\xi^2 + \tilde{C}_3)^{-2} \exp(iC_4), \quad n = 16, \quad (4.128)$$

$$U = (2/(3k))\xi^{-2} \exp(iC_4), \quad n = 13/3, \quad (4.129)$$

$$U = (2/k)\xi^{-2}(2\tilde{C}_3\xi^{1/3} + 3)(\tilde{C}_3\xi^{1/3} + 3)^{-2} \exp(iC_4), \quad n = 13/3, \quad (4.130)$$

with real constants \tilde{C}_3, C_4 .

D. Linearization of the translation-group ODE

When $C_1 = 0$ and $p = -1$, the level-set ODE (4.5) for $A(\xi)$ becomes linear,

$$A'' + (n-1)\xi^{-1}A' + \nu A + k = 0. \quad (4.131)$$

The general solution of this ODE splits into three cases.

For the case $\nu > 0$, we can use a change of variable $A = \xi^{1-n/2}B$, $z = \sqrt{\nu}\xi$ which transforms the ODE (4.131) to an inhomogeneous form of Bessel's equation

$$z^2 B'' + zB' + (z^2 - (\tfrac{1}{2}n - 1)^2)B + k\nu^{-(2+n)/4}z^{1+n/2} = 0. \quad (4.132)$$

This equation has the general solution

$$B = C_2 J_{|n-2|/2}(z) + C_3 Y_{|n-2|/2}(z) - k\nu^{-(2+n)/4}z^{|n-2|/2}, \quad (4.133)$$

where $J_\mu(z)$ and $Y_\mu(z)$ are the Bessel functions of the first and second kinds, respectively. Hence the general solution of ODE (4.131) for $A(\xi)$ with $\nu \neq 0$ is given by

$$A = \xi^{1-n/2}(C_2 J_{|n-2|/2}(\sqrt{\nu}\xi) + C_3 Y_{|n-2|/2}(\sqrt{\nu}\xi)) - k/\nu. \quad (4.134)$$

For the case $\nu < 0$, we can use a change of variable $A = \xi^{1-n/2}B$, $z = \sqrt{-\nu}\xi$ which transforms the ODE (4.131) to an inhomogeneous form of a modified Bessel equation

$$z^2 B'' + zB' - (z^2 + (\tfrac{1}{2}n - 1)^2)B + k(-\nu)^{-(2+n)/4}z^{1+n/2} = 0. \quad (4.135)$$

This equation has the general solution

$$B = C_2 I_{|n-2|/2}(z) + C_3 K_{|n-2|/2}(z) + k(-\nu)^{-(2+n)/4}z^{|n-2|/2}, \quad (4.136)$$

where $I_\mu(z)$ and $K_\mu(z)$ are the modified Bessel functions of the first and second kinds, respectively. Hence, the general solution of ODE (4.131) for $A(\xi)$ with $\nu \neq 0$ is given by

$$A = \xi^{1-n/2}(C_2 I_{|n-2|/2}(\sqrt{-\nu}\xi) + C_3 K_{|n-2|/2}(\sqrt{-\nu}\xi)) - k/\nu. \quad (4.137)$$

Since $C_1 = 0$, the first integral (3.31) for $\Phi(\xi)$ reduces to $\Phi' = 0$, which yields

$$\Phi = \text{const.} \quad (4.138)$$

For the case $\nu = 0$, the ODE (4.131) can be integrated directly to get

$$A = -(k/(2n))\xi^2 + C_3\xi^{2-n} + C_2, \quad n \neq 0, 2, \quad (4.139)$$

$$A = -(k/4)\xi^2 + C_3 \ln \xi + C_2, \quad n = 2, \quad (4.140)$$

$$A = (-k/2)\xi^2 \ln \xi + C_3\xi^2 + C_2, \quad n = 0. \quad (4.141)$$

From the first integral (3.31) for $\Phi(\xi)$, we again have

$$\Phi = \text{const.} \quad (4.142)$$

Hence we obtain the following result.

Proposition 12. For $p = -1$, the translation-group ODE (2.51) has solutions

$$U = (\xi^{1-n/2}(C_2 J_{|n-2|/2}(\sqrt{\nu}\xi) + C_3 Y_{|n-2|/2}(\sqrt{\nu}\xi)) - k/\nu) \exp(iC_4), \quad \nu > 0, \quad (4.143)$$

$$U = (\xi^{1-n/2}(C_2 I_{|n-2|/2}(\sqrt{-\nu}\xi) + C_3 K_{|n-2|/2}(\sqrt{-\nu}\xi)) - k/\nu) \exp(iC_4), \quad \nu < 0, \quad (4.144)$$

$$U = \left(-(k/(2n))\xi^2 + C_3\xi^{2-n} + C_2 \right) \exp(iC_4), \quad \nu = 0, \quad n \neq 0, 2 \quad (4.145)$$

$$U = \left(-(k/4)\xi^2 + C_3 \ln \xi + C_2 \right) \exp(iC_4), \quad \nu = 0, \quad n = 2, \quad (4.146)$$

$$U = \left((-k/2)\xi^2 \ln \xi + C_3\xi^2 + C_2 \right) \exp(iC_4), \quad \nu = 0, \quad n = 0, \quad (4.147)$$

with real constants C_2, C_3, C_4 .

V. SOLUTIONS TO THE OPTIMAL SCALING-GROUP ODE

Next the reduction of order methods from Sec. III will be applied to the scaling-group ODE (2.47) arising by the reduction of the radial Schrödinger equation under its optimal subgroup of point symmetries (2.42c).

To proceed, we use polar variables $U = A \exp(i\Phi)$ to convert this U(1)-invariant ODE (2.47) into a semilinear system of real ODEs

$$4\xi^2 A'' - 4\xi^2 A \Phi'^2 - 2(n-4-4/p)\xi A' + (1+4\mu\xi)A\Phi' \quad (5.1a)$$

$$- (\mu^2 - (4-2n)/p - 4/p^2)A + kA^{1+p} = 0,$$

$$4\xi^2 \Phi'' + 8\xi^2 A^{-1} A' \Phi' - (1+4\mu\xi)A^{-1} A' - 2(n-4-4/p)\xi \Phi' + \mu(n-2-4/p) = 0. \quad (5.1b)$$

As shown by the results stated in Proposition 1, the point symmetries of this system (5.1) consist of only phase rotations

$$Y_{\text{phas.}} = \partial/\partial\Phi. \quad (5.2)$$

In the case of the pseudo-conformal power $p = 4/n$, the ODEs in the polar system (5.1) are the respective Euler-Lagrange equations $\delta L/\delta A = 0$ and $A^{-1}\delta L/\delta\Phi = 0$ of the U(1)-invariant Lagrangian (2.48) expressed in polar variables,

$$L = -4\xi^2(A'^2 + A^2(\Phi' - (\frac{1}{8}\xi^{-2} + \frac{1}{2}\mu\xi^{-1}))^2) + (\frac{1}{16}\xi^{-2} + \frac{1}{2}\mu\xi^{-1} + 4\frac{p-1}{p^2})A^2 + \frac{2}{p+2}kA^{p+2}, \quad p = 4/n. \quad (5.3)$$

Invariance of the polar Lagrangian (5.3) under $Y_{\text{phas.}}$ produces a first integral (3.31), yielding

$$\Phi' = \frac{1}{8}\xi^{-2}(C_1 A^{-2} + 1 + 4\mu\xi), \quad p = 4/n. \quad (5.4)$$

The polar system (5.1) thereby reduces to a single real semilinear ODE

$$4\xi^2 A'' + 8\xi A' - \frac{1}{4}C_1^2 \xi^{-2} A^{-3} + \left(\frac{1}{16}\xi^{-2} + \frac{1}{2}\mu\xi^{-1} + 4\frac{p-1}{p^2}\right)A + kA^{1+p} = 0, \quad p = 4/n, \quad (5.5)$$

which is the Euler–Lagrange equation $\delta\tilde{L}/\delta A = 0$ of a modified Lagrangian (3.33), given by

$$L = -4\xi^2 A'^2 + \frac{1}{4}C_1^2 \xi^{-2} A^{-2} + \left(\frac{1}{16}\xi^{-2} + \frac{1}{2}\mu\xi^{-1} + 4\frac{p-1}{p^2}\right)A^2 + \frac{2}{p+2}kA^{p+2}, \quad p = 4/n. \quad (5.6)$$

Solutions of the ODE (5.5) for $A(\xi)$ represent the level set $C_1 = \text{const}$ of solutions $(A(\xi), \Phi(\xi))$ to the polar system (5.1), or equivalently the level set

$$C_1 = 2i\xi^2(U\bar{U}' - U'\bar{U}) - \left(\frac{1}{2} + 2\mu\xi\right)|U|^2 = \text{const} \quad (5.7)$$

of solutions $U(\xi)$ to the scaling-group ODE (2.47).

Note that the level-set ODE (5.5) will be linear if and only if $C_1 = 0$ and $p = -1$.

Proposition 13. In the nonlinear case $C_1 \neq 0$ or $p \neq -1$, the level-set ODE (5.5) admits no point symmetries.

Since both the level-set ODE (5.5) and the polar system (5.1) have no point symmetries, only the second reduction method (cf. Sec. III B) is applicable by using the linearization of ODE (5.5) which holds when $C_1 = 0, p = -1$.

A. Linearization of the scaling-group ODE

In the case $C_1 = 0$ and $p = -1$ ($n = -4$), the level-set ODE (5.5) for $A(\xi)$ becomes linear,

$$4\xi^2 A'' + 8\xi A' + \left(\frac{1}{16}\xi^{-2} + \frac{1}{2}\mu\xi^{-1} - 8\right)A + k = 0. \quad (5.8)$$

The general solution of this ODE is given by

$$\begin{aligned} A = & C_2 M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) + C_3 M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) \\ & + (ik/3) \left(M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) \int M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) d\xi \right. \\ & \left. - M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) \int M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) d\xi \right), \end{aligned} \quad (5.9)$$

where $M_{\lambda, \nu}(z)$ is the Whittaker function.

Since $C_1 = 0$, the first integral (3.31) for $\Phi(\xi)$ reduces to $\Phi' = \frac{1}{8}\xi^{-2} + \frac{1}{2}\mu\xi^{-1}$, which yields

$$\Phi = -\frac{1}{8}\xi^{-1} + \frac{1}{2}\mu \ln \xi + C_4. \quad (5.10)$$

Hence we obtain the following result.

Proposition 14. For $p = -1$, the scaling-group ODE (2.47) has a solution

$$\begin{aligned} U = & \left(C_2 M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) + C_3 M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) \right. \\ & + (ik/3) \left(M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) \int M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) d\xi \right. \\ & \left. \left. - M_{-i\mu/2, 3/2}(\tfrac{1}{4}i\xi^{-1}) \int M_{-i\mu/2, -3/2}(\tfrac{1}{4}i\xi^{-1}) d\xi \right) \right) \exp\left(-\frac{1}{8}\xi^{-1} + \frac{1}{2}i\mu \ln \xi + iC_4\right), \\ & n = -4, \end{aligned} \quad (5.11)$$

with real constants C_2, C_3, C_4 .

VI. SOLUTIONS TO THE OPTIMAL PSEUDO-CONFORMAL-GROUP ODE

The reduction of order methods from Sec. III will be considered last for the pseudo-conformal-group ODE (2.56) arising by the reduction of the radial Schrödinger equation under its optimal subgroup of point symmetries (2.42d).

We begin by using polar variables $U = A \exp(i\Phi)$ to convert this $U(1)$ -invariant ODE (2.56) into a semilinear system of real ODEs

$$4\xi^2 A'' - 4\xi^2 A \Phi'^2 + 8\xi A' + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} + n(1 - \frac{1}{4}n))A + kA^{1+4/n} = 0, \quad (6.1a)$$

$$\Phi'' + 2A^{-1}A'\Phi' + 2\xi^{-1}\Phi' = 0. \quad (6.1b)$$

The ODEs in this system (6.1) are the respective Euler–Lagrange equations $\delta L/\delta A = 0$ and $A^{-1}\delta L/\delta \Phi = 0$ of the $U(1)$ -invariant Lagrangian (2.57) expressed in polar variables,

$$L = -4\xi^2 A'^2 - 4\xi^2 A^2 \Phi'^2 + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} + n(1 - \frac{1}{4}n))A^2 + \frac{n}{n+2}kA^{2+4/n}. \quad (6.2)$$

As shown by the results stated in Proposition 1, the point symmetries of the polar system (6.1) consist of only phase rotations

$$Y_{\text{phas.}} = \partial/\partial \Phi. \quad (6.3)$$

Invariance of the polar Lagrangian (6.2) under $Y_{\text{phas.}}$ produces a first integral (3.31), yielding

$$\Phi' = \frac{1}{4}C_1 \xi^{-2} A^{-2}. \quad (6.4)$$

The polar system (6.1) thereby reduces to a single real semilinear ODE

$$4\xi^2 A'' + 8\xi A' - \frac{1}{4}C_1^2 \xi^{-2} A^{-3} + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} + n(1 - \frac{1}{4}n))A + kA^{1+4/n} = 0, \quad (6.5)$$

which is the Euler–Lagrange equation $\delta \tilde{L}/\delta A = 0$ of a modified Lagrangian (3.33), given by

$$\tilde{L} = -4\xi^2 A'^2 + \frac{1}{4}C_1^2 \xi^{-2} A^{-2} + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} + n(1 - \frac{1}{4}n))A^2 + \frac{n}{n+2}kA^{2+4/n}. \quad (6.6)$$

Solutions of the ODE (6.5) for $A(\xi)$ represent the level set $C_1 = \text{const}$ of solutions $(A(\xi), \Phi(\xi))$ to the polar system (6.1), or equivalently the level set

$$C_1 = 2i\xi^{-2}(U\bar{U}' - U'\bar{U}) = \text{const} \quad (6.7)$$

of solutions $U(\xi)$ to the pseudo-conformal-group ODE (2.56).

Note that the level-set ODE (6.5) will be linear if and only if $C_1 = 0$ and $p = -1$ ($n = -4$).

Proposition 15. In the nonlinear case $C_1 \neq 0$ or $p \neq -1$ ($n \neq -4$), the level-set ODE (6.5) admits no point symmetries.

Since both the level-set ODE (6.5) and the polar system (6.1) have no point symmetries, only the second reduction method (cf. Sec. IIIB) is applicable by using the linearization of ODE (6.5) which holds when $C_1 = 0$, $p = -1$ ($n = -4$).

A. Linearization of the pseudo-conformal-group ODE

In the case $C_1 = 0$ and $p = -1$ ($n = -4$), the level-set ODE (6.5) for $A(\xi)$ becomes linear,

$$4\xi^2 A'' + 8\xi A' + (\kappa \xi^{-1} - \frac{1}{4}\xi^{-2} - 8)A + k = 0. \quad (6.8)$$

The general solution of this ODE is given by

$$A = C_2 M_{\kappa/2, 3/2}(\frac{1}{2}\xi^{-1}) + C_3 M_{\kappa/2, -3/2}(\frac{1}{2}\xi^{-1}) + (k/6)(M_{\kappa/2, 3/2}(\frac{1}{2}\xi^{-1}) \int M_{\kappa/2, -3/2}(\frac{1}{2}\xi^{-1}) d\xi - M_{\kappa/2, -3/2}(\frac{1}{2}\xi^{-1}) \int M_{\kappa/2, 3/2}(\frac{1}{2}\xi^{-1}) d\xi), \quad (6.9)$$

where $M_{\lambda, \nu}(z)$ is the Whittaker function.

Since $C_1 = 0$, the first integral (3.31) for $\Phi(\xi)$ reduces to $\Phi' = 0$, which yields

$$\Phi = \text{const.} \quad (6.10)$$

Hence we obtain the following result.

Proposition 16. For $p = -1$, the pseudo-conformal-group ODE (2.56) has a solution

$$\begin{aligned} U = & \left(C_2 M_{\kappa/2, 3/2}(\tfrac{1}{2}\xi^{-1}) + C_3 M_{\kappa/2, -3/2}(\tfrac{1}{2}\xi^{-1}) \right. \\ & + (k/6) \left(M_{\kappa/2, 3/2}(\tfrac{1}{2}\xi^{-1}) \int M_{\kappa/2, -3/2}(\tfrac{1}{2}\xi^{-1}) d\xi \right. \\ & \left. \left. - M_{\kappa/2, -3/2}(\tfrac{1}{2}\xi^{-1}) \int M_{\kappa/2, 3/2}(\tfrac{1}{2}\xi^{-1}) d\xi \right) \right) \exp(iC_4), \quad n = -4, \end{aligned} \quad (6.11)$$

with real constants C_2, C_3, C_4 .

VII. GROUP-INVARIANT RADIAL SOLUTIONS

All of the group-invariant solutions derived in Secs. IV–VI through the optimal subgroups of point symmetries (2.42b)–(2.42d) will now be written out in the form $u = f(t, r)$ for the radial Schrödinger equation (1.2) and its alternative two-dimensional formulation (1.5). Wherever possible, solutions are merged into families that do not overlap.

A. Solutions from the optimal translation-symmetry subgroup

Theorem 5. The radial Schrödinger equations (1.2) and (1.5) have the following group-invariant solutions with respect to time translations (2.52):

$$u = (2(p(n-2)-2)/(kp^2))^{1/p} r^{-2/p} \exp(ic_1), \quad p \neq 2/(n-2), k/(p(n-2)-2) > 0, \quad (7.1)$$

$$\begin{aligned} u = & (\pm 8(p+2)/(kp^2))^{1/p} (c_2 r^2 \pm 1/c_2)^{-2/p} \exp(ic_1), \\ & p = 4/(n-2), \pm k(1-2/n) > 0, n \neq 2, \end{aligned} \quad (7.2)$$

$$\begin{aligned} u = & (-\tfrac{1}{8}kp^2(p+2))^{-1/p} (r + c_2 r^{3-n})^{-2/p} \exp(ic_1), \\ & p = 2(n-3)/(2-n), k/(n-2) < 0, n \neq 2, 3, \end{aligned} \quad (7.3)$$

$$u = (-\tfrac{1}{2}(k/n)r^2 + c_3 r^{2-n} + c_2) \exp(ic_1), \quad p = -1, n \neq 0, 2, \quad (7.4)$$

$$u = (-\tfrac{1}{4}kr^2 + c_3 \ln r + c_2) \exp(ic_1), \quad p = -1, n = 2, \quad (7.5)$$

$$u = (96/k)(r^2 + c_2)(2r^2 + c_2)^{-2} \exp(ic_1), \quad p = 1, n = 16, \quad (7.6)$$

$$u = (-\tfrac{1}{2}kr^2 \ln r + c_3 r^2 + c_2) \exp(ic_1), \quad p = -1, n = 0, \quad (7.7)$$

$$\begin{aligned} u = & (4k/3)^{1/4} r^2 (r^{-2} - c_2/k)^{3/4} \exp(ic_1 - i\sqrt{3}((k/c_2)r^{-2} - 1)^{-1/2}), \\ & p = -4, n = 0, k > 0, \end{aligned} \quad (7.8)$$

$$\begin{aligned} u = & (-12k)^{1/4} (c_2/k - r^{2/3})^{3/4} \exp(ic_1 + i\sqrt{3}((k/c_2)r^{2/3} - 1)^{-1/2}), \\ & p = -4, n = 4/3, k < 0, \end{aligned} \quad (7.9)$$

$$u = ((3k/c_3^2)r^{4/3} - 4c_3^2(r^{1/3} + c_2r^{2/3})^2)^{1/4} \times \exp\left(\mathrm{i}c_1 + \frac{1}{2}\mathrm{i}\arcsin\left(\left(1 - \frac{4}{3}(c_3^4/k)(r^{-1/3} + c_2)^2\right)^{1/2}\right)\right), \quad (7.10)$$

$$p = -8, n = 5/3, k > 0,$$

$$u = (2/k)r^{-2}(2c_2r^{1/3} + 3)(c_2r^{1/3} + 3)^{-2}\exp(\mathrm{i}c_1), \quad p = 1, n = 13/3, \quad (7.11)$$

$$u = (-2c_3r + (\pm(k - c_3^2))^{1/2}(c_2r^2 \mp 1/c_2))^{1/2}\exp(\mathrm{i}c_1), \quad p = -4, n = 1, c_3^2 \neq k, \quad (7.12)$$

$$u = (-2c_4r + (k - c_3^2 - c_4^2)^{1/2}(c_2r^2 - 1/c_2))^{1/2} \times \exp\left(\mathrm{i}c_1 - \frac{1}{2}\mathrm{i}(k/c_3^2 - 1)^{-1/2}\operatorname{arcsinh}\left(\frac{2(k - c_3^2 - c_4^2)^{1/2}r + c_4(c_2r^2 - 1/c_2)}{-2c_4r + (k - c_3^2 - c_4^2)^{1/2}(c_2r^2 - 1/c_2)}\right)\right),$$

$$p = -4, n = 1, c_3^2 + c_4^2 < k, \quad (7.13)$$

$$u = (-2c_4r + (c_3^2 + c_4^2 - k)^{1/2}(c_2r^2 + 1/c_2))^{1/2} \times \exp\left(\mathrm{i}c_1 - \frac{1}{2}\mathrm{i}(k/c_3^2 - 1)^{-1/2}\operatorname{arccosh}\left(\frac{2(c_3^2 + c_4^2 - k)^{1/2}r - c_4(c_2r^2 + 1/c_2)}{2c_4r - (c_3^2 + c_4^2 - k)^{1/2}(c_2r^2 + 1/c_2)}\right)\right),$$

$$p = -4, n = 1, c_3^2 < k < c_3^2 + c_4^2, \quad (7.14)$$

$$u = (-2c_4r + (c_3^2 + c_4^2 - k)^{1/2}(c_2r^2 + 1/c_2))^{1/2} \times \exp\left(\mathrm{i}c_1 - \frac{1}{2}\mathrm{i}(1 - k/c_3^2)^{-1/2}\arcsin\left(\frac{2(c_3^2 + c_4^2 - k)^{1/2}r - c_4(c_2r^2 + 1/c_2)}{-2c_4r + (c_3^2 + c_4^2 - k)^{1/2}(c_2r^2 + 1/c_2)}\right)\right),$$

$$p = -4, n = 1, c_3^2 > k, \quad (7.15)$$

$$u = (-2c_3r + c_2r^2)^{1/2}\exp(\mathrm{i}c_1 \pm \mathrm{i}\frac{1}{2}(k/c_3^2 - 1)^{1/2}\ln|c_2 - 2c_3/r|), \quad p = -4, n = 1, \quad (7.16)$$

$$u = (-2c_3r + c_2)^{1/2}\exp(\mathrm{i}c_1 \pm \mathrm{i}\frac{1}{2}(k/c_3^2 - 1)^{1/2}\ln|c_2 - 2c_3r|), \quad p = -4, n = 1, \quad (7.17)$$

$$u = (c_3/c_2)^{1/2}(c_2r - 1)\exp\left(\mathrm{i}c_1 - \frac{1}{2}\mathrm{i}(k^{1/2}/c_3)\left(\frac{c_2r + 1}{c_2r - 1}\right)\right), \quad p = -4, n = 1, k > 0, \quad (7.18)$$

$$u = c_2r\exp(\mathrm{i}c_1 - \mathrm{i}(k^{1/2}/c_2^2)r^{-1}), \quad p = -4, n = 1, k > 0, \quad (7.19)$$

$$u = c_2\exp(\mathrm{i}c_1 - \mathrm{i}(k^{1/2}/c_2^2)r), \quad p = -4, n = 1, k > 0, \quad (7.20)$$

$$u = (((c_3^2 - k)/c_4)r^2 + c_4(1 + c_2r)^2)^{1/2} \times \exp\left(\mathrm{i}c_1 - \mathrm{i}c_3(k - c_3^2)^{-1/2}\operatorname{arcsinh}\left(\left(-1 + (c_4^2/(k - c_3^2))(r^{-1} + c_2)^2\right)^{-1/2}\right)\right),$$

$$p = -4, n = 1, c_3^2 < k, c_4 > 0, \quad (7.21)$$

$$u = ((c_3^2 - k)/c_4)r^2 + c_4(1 + c_2r)^2)^{1/2} \times \exp\left(\mathrm{i}c_1 - \mathrm{i}c_3(c_3^2 - k)^{-1/2} \arcsin\left((1 + (c_4^2/(c_3^2 - k))(r^{-1} + c_2)^2)^{-1/2}\right)\right), \quad (7.22)$$

$$p = -4, n = 1, c_3^2 > k, c_4 > 0,$$

$$u = ((c_3^2 - k)/c_4)r^2 + c_4(1 + c_2r)^2)^{1/2} \times \exp\left(\mathrm{i}c_1 - \mathrm{i}c_3(k - c_3^2)^{-1/2} \operatorname{arccosh}\left((1 + (c_4^2/(c_3^2 - k))(r^{-1} + c_2)^2)^{-1/2}\right)\right), \quad (7.23)$$

$$p = -4, n = 1, c_3^2 < k, c_4 < 0,$$

$$u = (4(k - c_3^2))^{1/4} (c_2r^2 \pm r)^{1/2} \exp\left(\mathrm{i}c_1 + \frac{1}{2}\mathrm{i}c_3(k - c_3^2)^{-1/2} \ln|c_2 \pm r^{-1}|\right), \quad (7.24)$$

$$p = -4, n = 1, c_3^2 < k,$$

$$u = c_3(1 + c_2r) \exp\left(\mathrm{i}c_1 \mp \mathrm{i}(k^{1/2}/c_3^2)(r^{-1} + c_2)^{-1}\right), \quad p = -4, n = 1. \quad (7.25)$$

Theorem 6. The radial Schrödinger equations (1.2) and (1.5) have the following group-invariant solutions with respect to time translations combined with phase rotations (2.44b):

$$u = \left(r^{1-n/2}(c_2J_{|n-2|/2}(\sqrt{v}r) + c_3Y_{|n-2|/2}(\sqrt{v}r)) - k/v\right) \exp(\mathrm{i}c_1 + \mathrm{i}vt), \quad p = -1, \quad (7.26)$$

$$u = \left(r^{1-n/2}(c_2I_{|n-2|/2}(\sqrt{-v}r) + c_3K_{|n-2|/2}(\sqrt{-v}r)) - k/v\right) \exp(\mathrm{i}c_1 + \mathrm{i}vt), \quad p = -1. \quad (7.27)$$

B. Solutions from the optimal scaling-symmetry and pseudo-conformal-symmetry subgroups

Theorem 7. The radial Schrödinger equations (1.2) and (1.5) have the following group-invariant solutions with respect to scalings combined with phase rotations (2.44c):

$$u = \frac{1}{3}\mathrm{i}kr^2 \left(M_{-\mathrm{i}\mu/2, -3/2}(\frac{1}{4}\mathrm{i}r^2/t) \int_{c_2}^{t/r^2} M_{-\mathrm{i}\mu/2, 3/2}(\frac{1}{4}\mathrm{i}\xi^{-1}) d\xi \right. \\ \left. - M_{-\mathrm{i}\mu/2, 3/2}(\frac{1}{4}\mathrm{i}r^2/t) \int_{c_3}^{t/r^2} M_{-\mathrm{i}\mu/2, -3/2}(\frac{1}{4}\mathrm{i}\xi^{-1}) d\xi \right) \exp\left(\mathrm{i}c_1 - \frac{1}{8}\mathrm{i}(r^2/t) + \frac{1}{2}\mathrm{i}\mu \ln t\right), \\ p = -1, n = -4. \quad (7.28)$$

Theorem 8. The radial Schrödinger equations (1.2) and (1.5) have the following group-invariant solutions with respect to time translations and inversions combined with phase rotations (2.44d):

$$u = \frac{1}{6}kr^2 \left(M_{\kappa/2, 3/2}(\frac{1}{2}r^2/(1+t^2)) \int_{c_2}^{(1+t^2)/r^2} M_{\kappa/2, -3/2}(\frac{1}{2}\xi^{-1}) d\xi \right. \\ \left. - M_{\kappa/2, -3/2}(\frac{1}{2}r^2/(1+t^2)) \int_{c_3}^{(1+t^2)/r^2} M_{\kappa/2, 3/2}(\frac{1}{2}\xi^{-1}) d\xi \right) \\ \times \exp\left(\mathrm{i}c_1 - \mathrm{i}\kappa \arctan(1/t) - \frac{1}{4}\mathrm{i}r^2t/(1+t^2)\right), \\ p = -1, n = -4. \quad (7.29)$$

C. Solutions from the full symmetry group

From Theorem 3, the full group of point symmetries admitted by the radial Schrödinger equation (1.2) and its alternative two-dimensional formulation (1.5) can be applied to each of the solutions $u = f(t, r)$ listed in Theorems 5–8. Phase rotations (2.32a) and scalings (2.32c) change only the constants appearing in these solutions, while time translations (2.32b) at most shift t by a new constant. In contrast, inversions (2.32d) produce additional new solutions, which are listed in the next theorem.

Theorem 9. *In the case of the pseudo-conformal power $p = 4/n$, the radial Schrödinger equations (1.2) and (1.5) have the following additional group-invariant solutions:*

$$u = (n(n-4)/(4k))^{n/4} r^{-n/2} \exp(ic_1 - ic_2 r^2/(4(1+c_2t))),$$

$$n \neq 4, k/(n(n-4)) > 0,$$
(7.30)

$$u = (-n^3/(2k(2n+4)))^{n/4} r^{-n/2} (1 + c_2(r/(1+c_3t))^{2-4/p})^{-2/p} \exp(ic_1 - ic_3 r^2/(4(1+c_3t))),$$

$$n^2 - n + 4 = 0,$$
(7.31)

$$u = (\frac{1}{8}kr^2 + c_3r^6(1+c_4t)^{-4} + c_2(1+c_4t)^2) \exp(ic_1 - ic_4r^2/(1+c_4t)),$$

$$p = -1, n = -4,$$
(7.32)

$$u = \left(r^3(1+c_4t)^{-1} (c_2J_3(\sqrt{v}r/(1+c_4t)) + c_3Y_3(\sqrt{v}r/(1+c_4t))) - (k/v)(1+c_4t)^2 \right)$$

$$\times \exp(ic_1 + ivt/(1+c_4t) - ic_4r^2/(4(1+c_4t))),$$
(7.33)

$$p = -1, n = -4, v > 0,$$

$$u = \left(r^3(1+c_4t)^{-1} (c_2I_3(\sqrt{-v}r/(1+c_4t)) + c_3K_3(\sqrt{-v}r/(1+c_4t))) - (k/v)(1+c_4t)^2 \right)$$

$$\times \exp(ic_1 + ivt/(1+c_4t) - ic_4r^2/(4(1+c_4t))),$$

$$p = -1, n = -4, v < 0,$$
(7.34)

$$u = \frac{1}{3}ikr^2 \left(M_{-i\mu/2, -3/2}(\frac{1}{4}ir^2/(t(1+c_4t))) \int_{c_2}^{t(1+c_4t)/r^2} M_{-i\mu/2, 3/2}(\frac{1}{4}i\xi^{-1}) d\xi \right.$$

$$\left. - M_{-i\mu/2, 3/2}(\frac{1}{4}ir^2/(t(1+c_4t))) \int_{c_3}^{t(1+c_4t)/r^2} M_{-i\mu/2, -3/2}(\frac{1}{4}i\xi^{-1}) d\xi \right)$$

$$\times \exp(ic_1 - \frac{1}{2}i\mu \ln(c_4 + 1/t) - \frac{1}{8}ir^2(1+2c_4t)/(t(1+c_4t))),$$

$$p = -1, n = -4,$$

$$u = \frac{1}{6}kr^2 \left(M_{\kappa/2, 3/2}(\frac{1}{2}r^2/(t^2 + (1+c_4t)^2)) \int_{c_2}^{(t^2+(1+c_4t)^2)/r^2} M_{\kappa/2, -3/2}(\frac{1}{2}\xi^{-1}) d\xi \right.$$

$$\left. - M_{\kappa/2, -3/2}(\frac{1}{2}r^2/(t^2 + (1+c_4t)^2)) \int_{c_3}^{(t^2+(1+c_4t)^2)/r^2} M_{\kappa/2, 3/2}(\frac{1}{2}\xi^{-1}) d\xi \right)$$

$$\times \exp(ic_1 - i\kappa \arctan(c_4 + 1/t) - \frac{1}{4}ir^2(c_4 + t + c_4^2t)/(t^2 + (1+c_4t)^2)),$$

$$p = -1, n = -4.$$
(7.36)

D. Analytical behaviour of group-invariant solutions

The radial solutions listed in Theorems 5–9 exhibit several types of interesting behaviour describing (1) standing waves; (2) static and dynamic monopoles; (3) static “bright solitons”; (4) static and dynamic “dark solitons”.

A function $u(t, r)$ describes a *radial monopole* if it is smooth on $0 < r < \infty$ such that $|u| \rightarrow 0$ as $r \rightarrow \infty$ and $|u| \rightarrow \infty$ as $r \rightarrow 0$. Solution (7.1) for $p > 0$ is a n -dimensional static monopole; solution (7.6) for $c_2 = 0$ is a 16-dimensional static monopole; solutions (7.1) for $n \neq 1, 2, \dots$, (7.3) for $2 < n < 3$, and (7.11) are planar static monopoles. Solution (7.30) is a n -dimensional monopole with a dynamic phase; solution (7.31) is a planar dynamic monopole.

A function $u(t, r)$ describes a *radial standing wave* if it is smooth on $0 \leq r < \infty$ such that $u = U(r)\exp(i\omega t)$ with $\omega \neq 0$ and $|U|$ bounded as $r \rightarrow \infty$. Solution (7.26) and solution (7.27) in the case $c_2 = 0$ are n -dimensional standing waves.

A function $u(t, r)$ describes a “*bright radial soliton*” or a “*dark radial soliton*” if it is smooth on $0 \leq r < \infty$ such that $|u| \rightarrow 0$ as $r \rightarrow \infty$ or $|u| \rightarrow A \neq 0$ as $r \rightarrow \infty$, respectively. Solution (7.2) for $n > 2$ in the “+” case is a n -dimensional static bright soliton; solution (7.6) in the case $c_2 = 0$ is a 16-dimensional static bright soliton. Solution (7.34) in the case $c_2 = 0$ is a planar dynamic dark soliton exhibiting blow-up.

More details about these physically interesting radial solutions, including their L^2 norms and conserved energies, will be discussed in a separate paper.¹³ The remaining radial solutions found in Theorems 5–9 all have unphysical behaviour, in particular $|u|$ is unbounded as $r \rightarrow \infty$.

VIII. CONCLUDING REMARKS

In this paper, all explicit group-invariant solutions given by elementary functions have been derived (cf. Theorems 5–9) for the class of semilinear radial Schrödinger equations (1.2) with a power nonlinearity $p \neq 0$ in multi-dimensions $n \neq 1$. Among these solutions $u(t, r)$, a few describe n -dimensional radial standing waves, radial monopoles, and static radial “bright solitons”, which have some physical interest.

Several solutions exist, surprisingly, only for non-integer values of n . In such cases the radial Schrödinger equation (1.2) is shown to have an alternative interpretation as a planar (i.e., 2-dimensional) radial equation (1.5) containing an extra modulation term mu_r/r that describes a point-source disturbance at the origin $r = 0$, with $m = 2 - n$. Some of these planar solutions are physically interesting dynamic radial monopoles and dynamic radial “dark solitons”.

However, no n -dimensional radial solutions are obtained in the analytically important cases $p \geq 4/n$ relevant for blow-up behaviour when $n \geq 2$. In particular, it is rigorously known^{1,2} that some radial solutions exhibit a finite time blow-up such that $|u(t, r)| \rightarrow \infty$ as $t \rightarrow T < \infty$ (with the energy and L^2 norm of $u(t, r)$ being finite).

For the critical case $p = 4/n$, a special class of radial blow-up solutions can be shown to have an exact group-invariant form¹

$$u(t, r) = (T - t)^{-n/2} U(\xi) \exp(i(\omega + r^2/4)/(T - t)), \quad \xi = r/(T - t), \quad \omega \neq 0, \quad (8.1)$$

where $U(\xi)$ satisfies the complex nonlinear 2nd order ODE

$$U'' + (n - 1)\xi^{-1}U' + \omega U + k|U|^{4/n}U = 0 \quad (8.2)$$

which is given by reduction of the radial Schrödinger equation (1.2) under a pseudo-conformal symmetry group generated by

$$T^2 X_{\text{trans.}} - T X_{\text{scal.}} + X_{\text{inver.}} + \omega X_{\text{phas.}} \quad (8.3)$$

Such pseudo-conformal blow-up solutions (8.1) are related by a certain symmetry transformation to standing wave solutions (2.50) and (2.51), which arise from the optimal translation group (2.44b) given by the generator (2.42b). Specifically, if $u = f(t, r)$ has a standing-wave form with frequency $\nu \neq 0$ when $p = 4/n$, then

$$u = f(t/(T - t), r/(T - t))(T - t)^{-2/p} \exp(ir^2/(4(T - t))) \quad (8.4)$$

(modulo a constant phase rotation) has a blow-up form with $\omega = T\nu \neq 0$, where

$$t \rightarrow t/(1 - t/T), \quad r \rightarrow r/(T - t), \quad u \rightarrow (T - t)^{2/p} \exp(-ir^2/(4(T - t)))u \quad (8.5)$$

is an inversion transformation combined with a scaling transformation acting on (t, r, u) . Hence the ODEs (2.51) and (8.2) are equivalent up to at most a point transformation on (ξ, U, \bar{U}) . From the results in Propositions 1(ii) and 2, we conclude that, since $\omega \neq 0$, the blow-up ODE (8.2) has too few point symmetries to allow it to be reduced to quadratures by means of first integrals (except in the case $p \neq -1, n = 4/p = -4$ when special solutions can be found in terms of Bessel functions). The only obvious explicit solution, by inspection, is $U = (-\omega/k)^{n/4}$, but this solution has infinite energy and infinite L^2 norm.

In the supercritical case $p > 4/n$, numerical evidence¹ suggests that a general class of blow-up solutions for the radial Schrödinger equation (1.2) asymptotically approach an exact similarity form

$$u(t, r) = (T - t)^{-1/p} U(\xi) \exp(i\omega \ln((T - t)/T)), \quad \xi = r/\sqrt{T - t}, \quad \omega \neq 0, \quad (8.6)$$

where $U(\xi)$ satisfies a more complicated complex nonlinear 2nd order ODE

$$U'' + ((n - 1)\xi^{-1} - \frac{1}{2}i\xi)U' - (\omega + i/p)U + k|U|^p U = 0, \quad (8.7)$$

which is given by reduction with respect to a scaling symmetry group generated by

$$TX_{\text{trans.}} - \frac{1}{2}X_{\text{scal.}} - \omega X_{\text{phas.}} \quad (8.8)$$

Modulo a time translation symmetry transformation $t \rightarrow t + T$, these blow-up solutions (8.6) are the same as the similarity solutions (2.46) and (2.47), which arise from the optimal scaling group (2.44c) given by the generator (2.42c). Since the ODEs (2.47) and (8.7) are thereby equivalent up to at most a point transformation on (ξ, U, \bar{U}) , the results in Propositions 1(i) and 13 show that the blow-up ODE (8.7) has too few point symmetries to allow it to be reduced to quadratures by means of first integrals (except, again, in the case $p \neq -1, n = 4/p = -4$ when special solutions can be found in terms of Whittaker functions).

Consequently, symmetry reduction methods are unable to yield any explicit n -dimensional radial blow-up solutions (8.1) and (8.6). To look for such solutions, we plan to apply the method of group-foliation reduction,¹⁴ which has been successfully used in previous work^{15,16} to obtain blow-up and dispersive radial solutions to semilinear wave equations and semilinear heat conduction equations with power nonlinearities in multi-dimensions.

ACKNOWLEDGMENTS

S. Anco is supported by a NSERC research grant. W. Feng is indebted to the China Scholarship council for financial support to work as a visiting scholar at Brock University and thanks the Department of Mathematics for support during the period when this paper was written. The referees are thanked for valuable comments which have improved this paper.

- ¹ C. Sulem and P.-L. Sulem, *The Nonlinear Schrödinger Equation*, Applied Mathematical Science (Springer, New York, 1999), Vol. 139.
- ² T. Cazenave, *Semilinear Schrödinger Equations*, Courant Lecture Notes 10 (American Mathematical Society, Providence, 2003).
- ³ A. D. Polyanin and V. F. Zaitsev, *Handbook of Nonlinear Partial Differential Equations*, 2nd ed. (CRC/Chapman and Hall, 2011).
- ⁴ P. J. Olver, *Applications of Lie Groups to Differential Equations* (Springer, New York, 1986).
- ⁵ G. Bluman and S. C. Anco, *Symmetry and Integration Methods for Differential Equations*, Applied Mathematical Science (Springer, New York, 2002), Vol. 154.
- ⁶ A. G. Nikitin and R. O. Popovych, *Ukr. Math. J.* **53**(8), 1255–1265 (2001).
- ⁷ W. I. Fushchich and N. I. Serov, *J. Phys. A* **20**, L929–L933 (1987).
- ⁸ L. Gagnon, and P. Winternitz, *J. Phys. A* **21**, 1493–1511 (1988); **22**, 469–497 (1989); **22**, 499–509 (1989).
- ⁹ L. Gagnon and P. Winternitz, *Phys. Rev. A* **39**, 296–306 (1989).
- ¹⁰ R. O. Popovych, M. Kunzinger, and H. Eshraghi, *Acta Appl. Math.* **109**, 315–359 (2010).
- ¹¹ J. Patera and P. Winternitz, *J. Math. Phys.* **18**, 1449–1455 (1977).
- ¹² S. C. Anco and W. Feng, in preparation (unpublished).
- ¹³ S. C. Anco, in preparation (unpublished).
- ¹⁴ S. C. Anco, S. Ali, and T. Wolf, *Symmetry, Integr. Geom.: Methods Appl.* **7**, 066 (2011).
- ¹⁵ S. C. Anco and S. Liu, *J. Math. Anal. Appl.* **297**, 317–342 (2004).
- ¹⁶ S. C. Anco, S. Ali, and T. Wolf, *J. Math. Anal. Appl.* **379**, 748–763 (2011).